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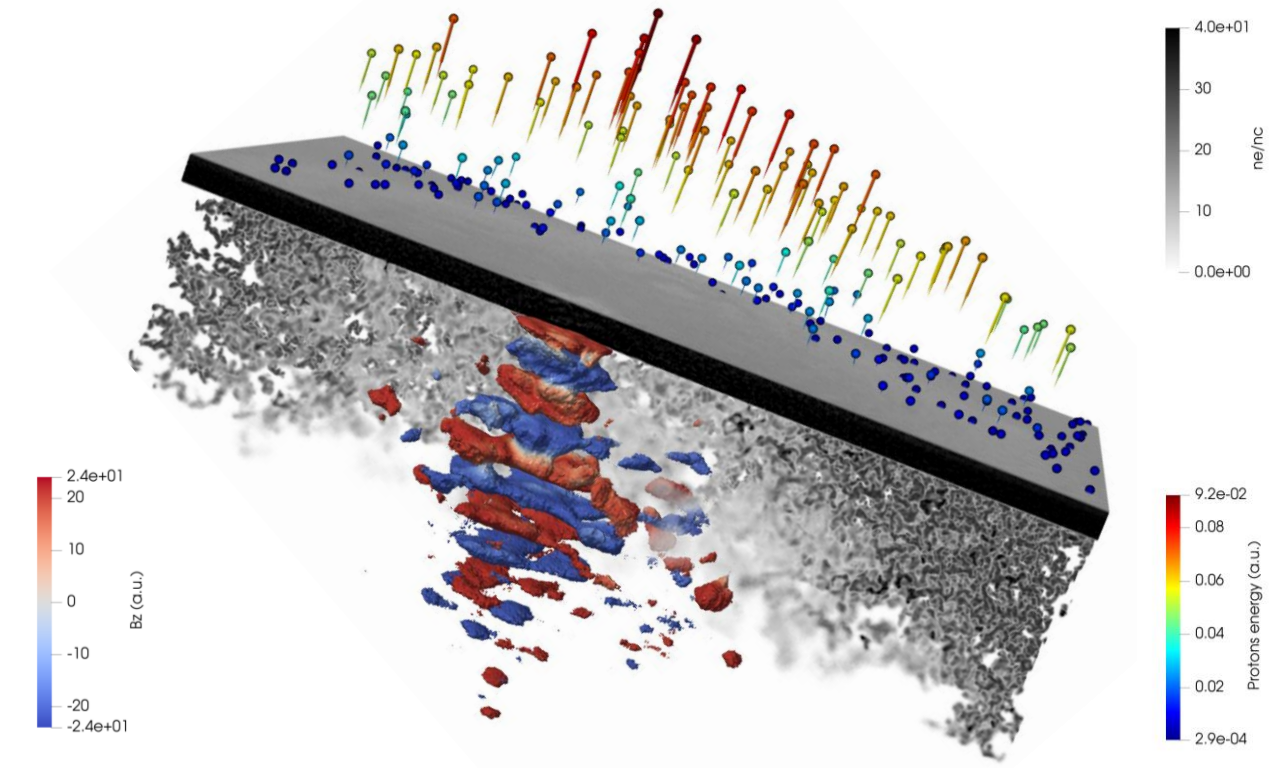


Department of Energy



ERC-2014-CoG No. 647554

ENSURE

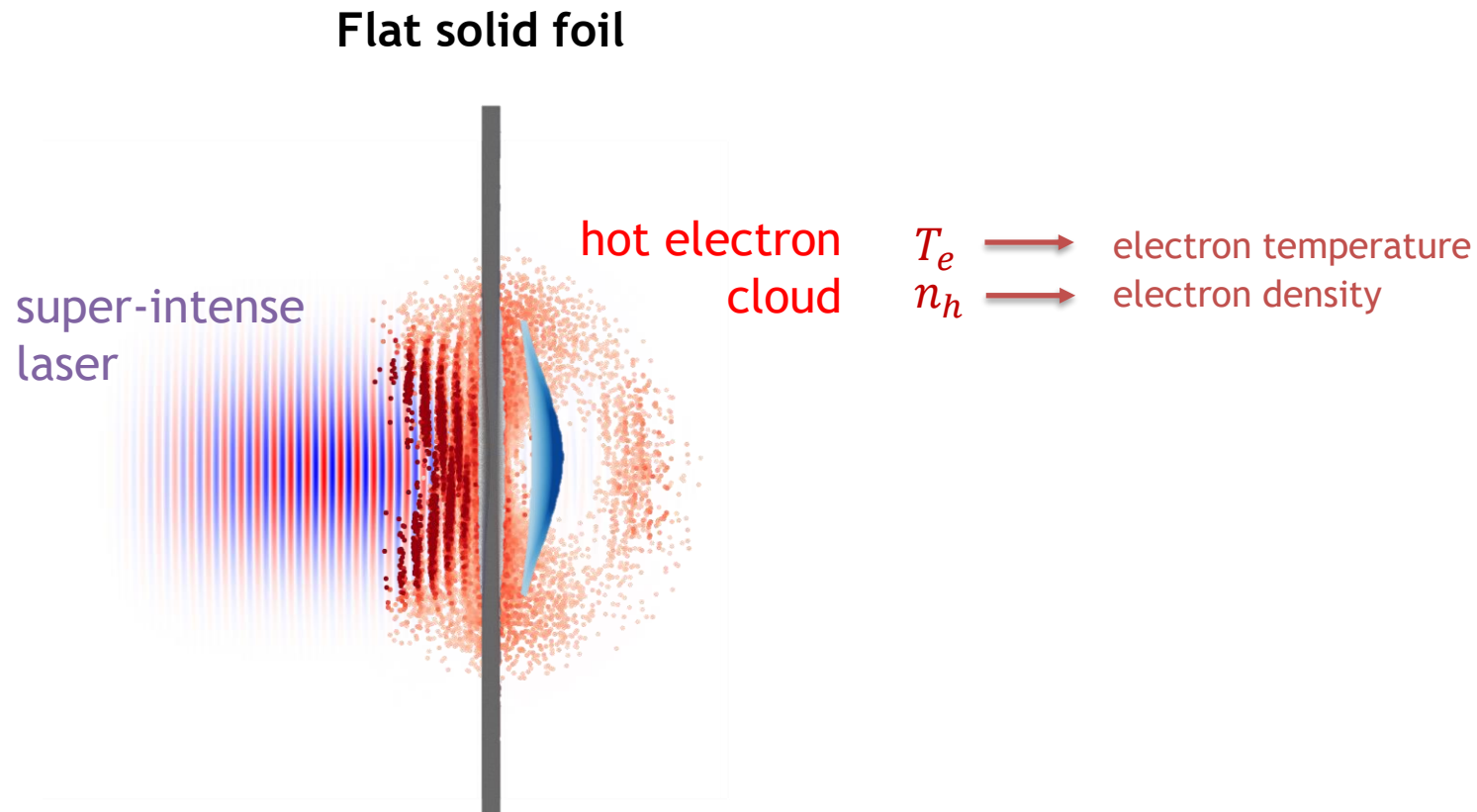


A theoretical model of laser-driven ion acceleration from near-critical double-layer targets

Andrea Pazzaglia, Luca Fedeli, Arianna Formenti, Alessandro Maffini, Matteo Passoni

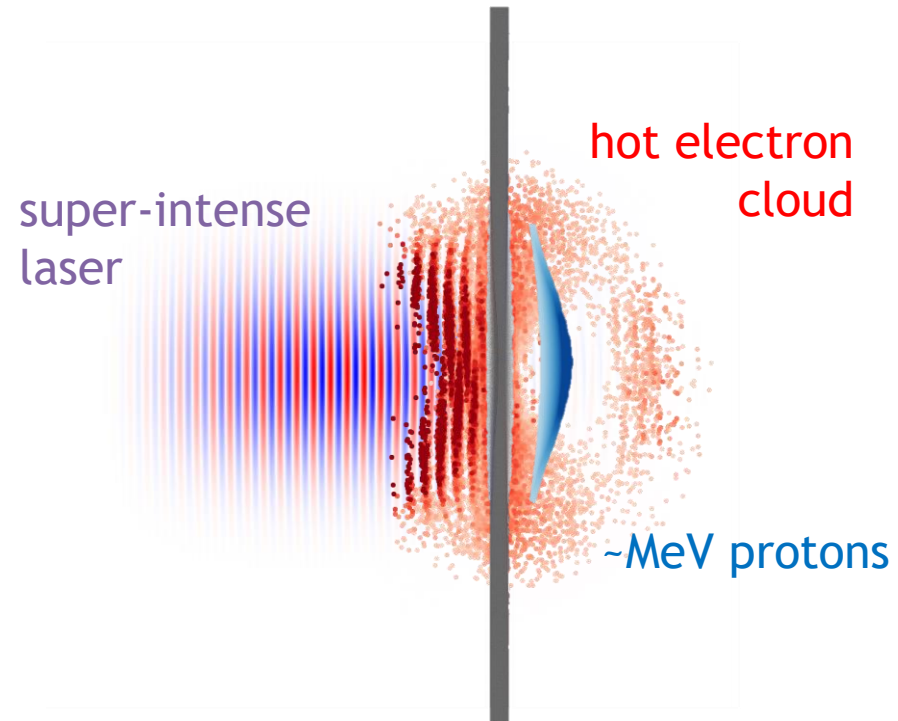
Department of Energy, Politecnico di Milano, Italy

Conventional laser-driven ion acceleration



Conventional laser-driven ion acceleration

Flat solid foil



hot electron
cloud

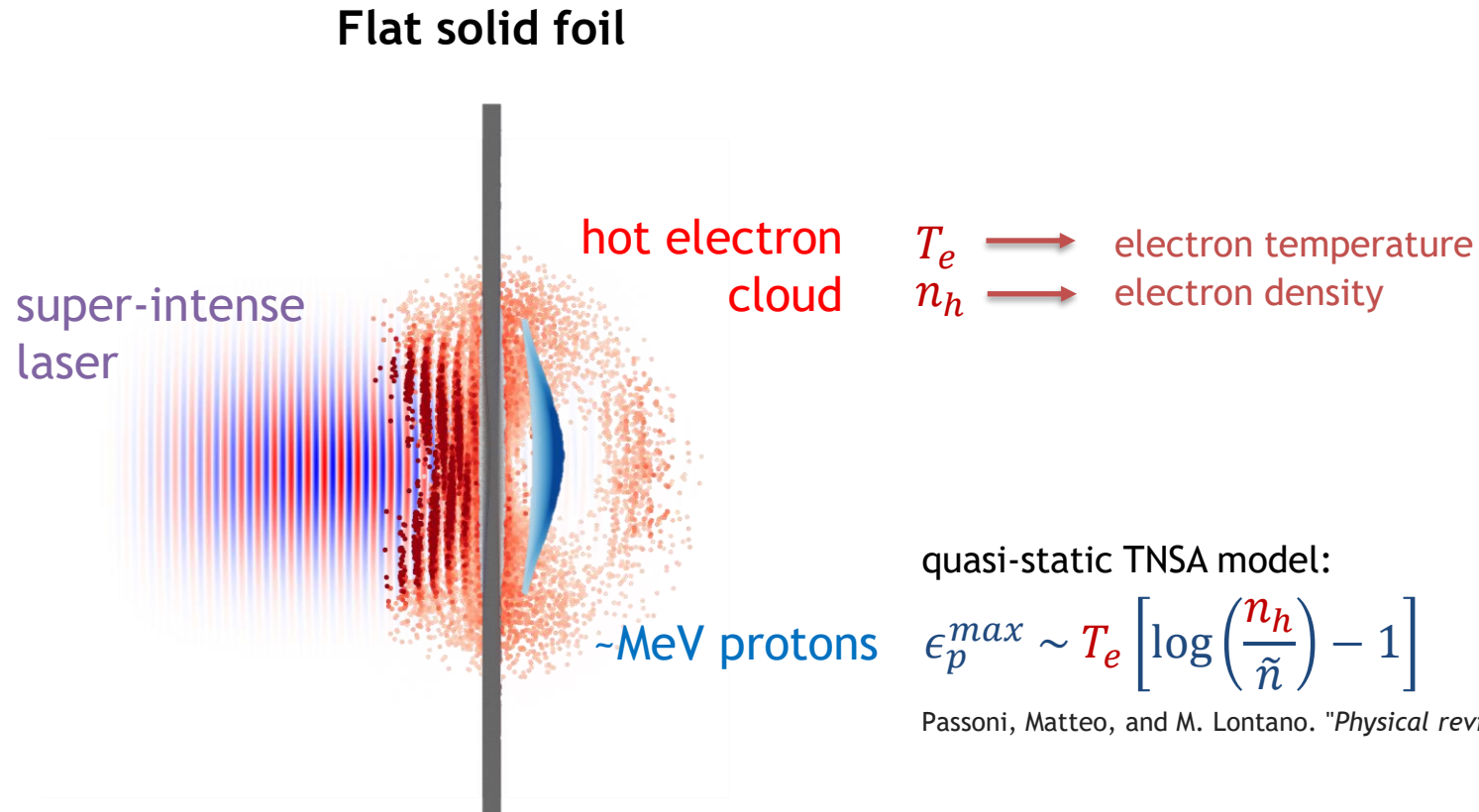
T_e → electron temperature
 n_h → electron density

quasi-static TNSA model:

$$\epsilon_p^{max} \sim T_e \left[\log \left(\frac{n_h}{\tilde{n}} \right) - 1 \right]$$

Passoni, Matteo, and M. Lontano. "Physical review letters 101.11 (2008): 115001.

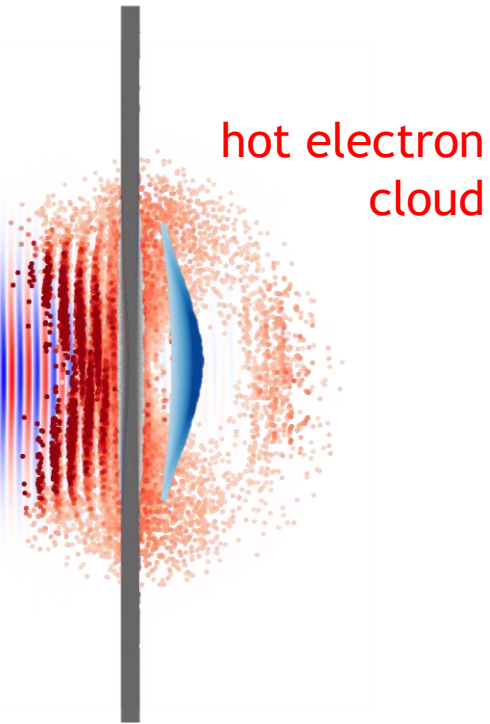
Conventional laser-driven ion acceleration



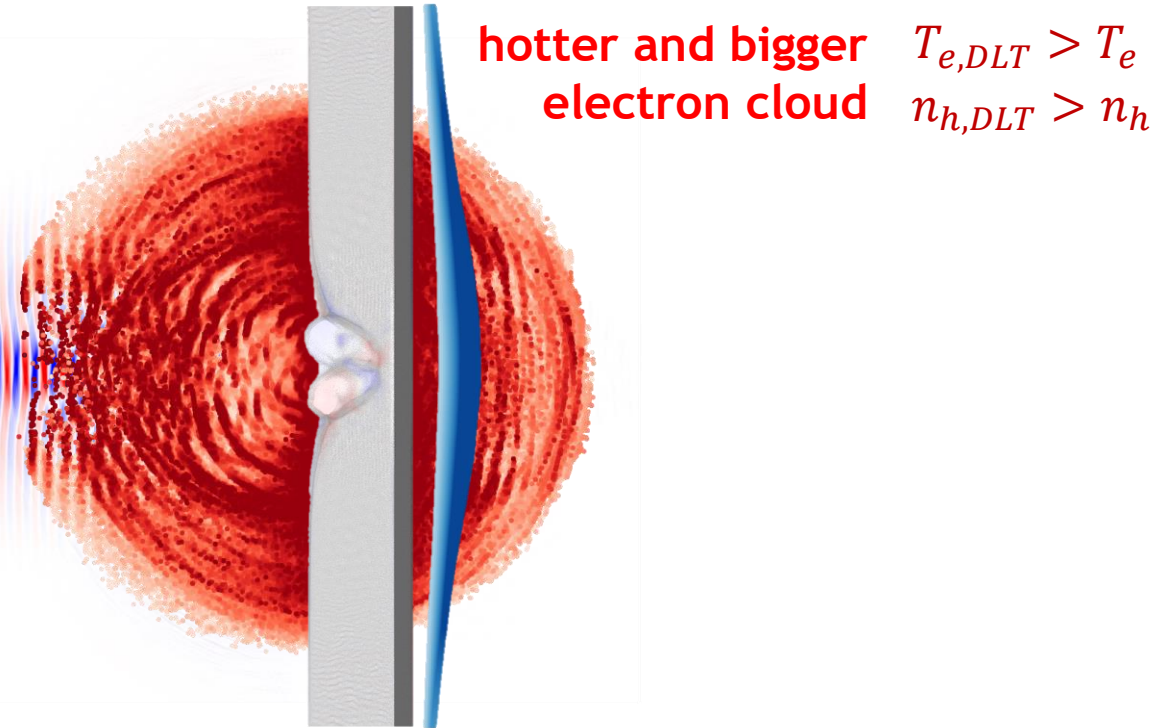
low electron conversion efficiency → low maximum proton energy

Acceleration with the Double-Layer Target (DLT)

Flat solid foil

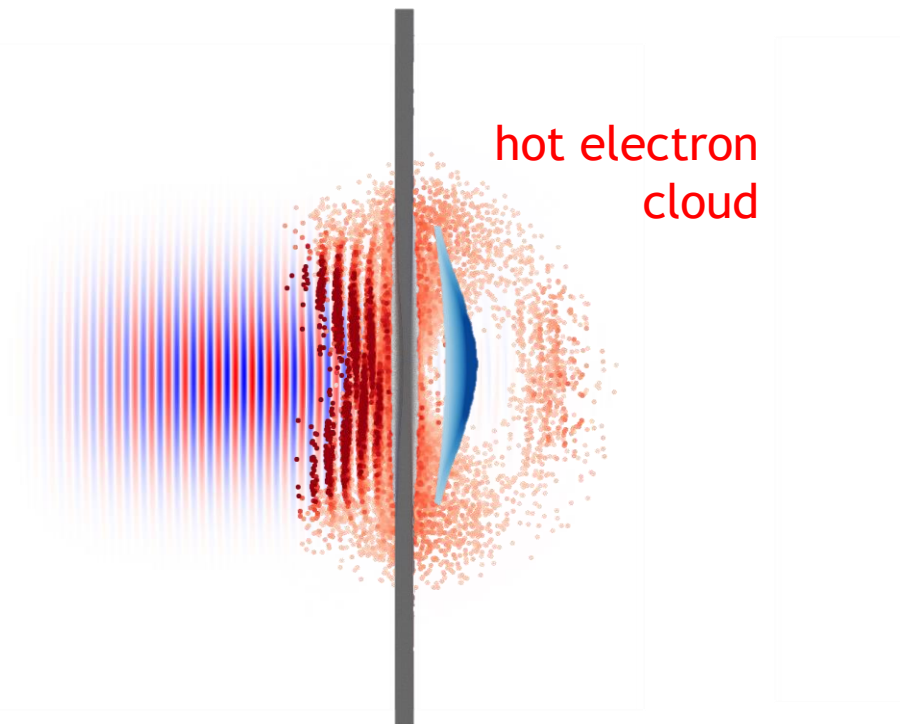


Near critical density layer

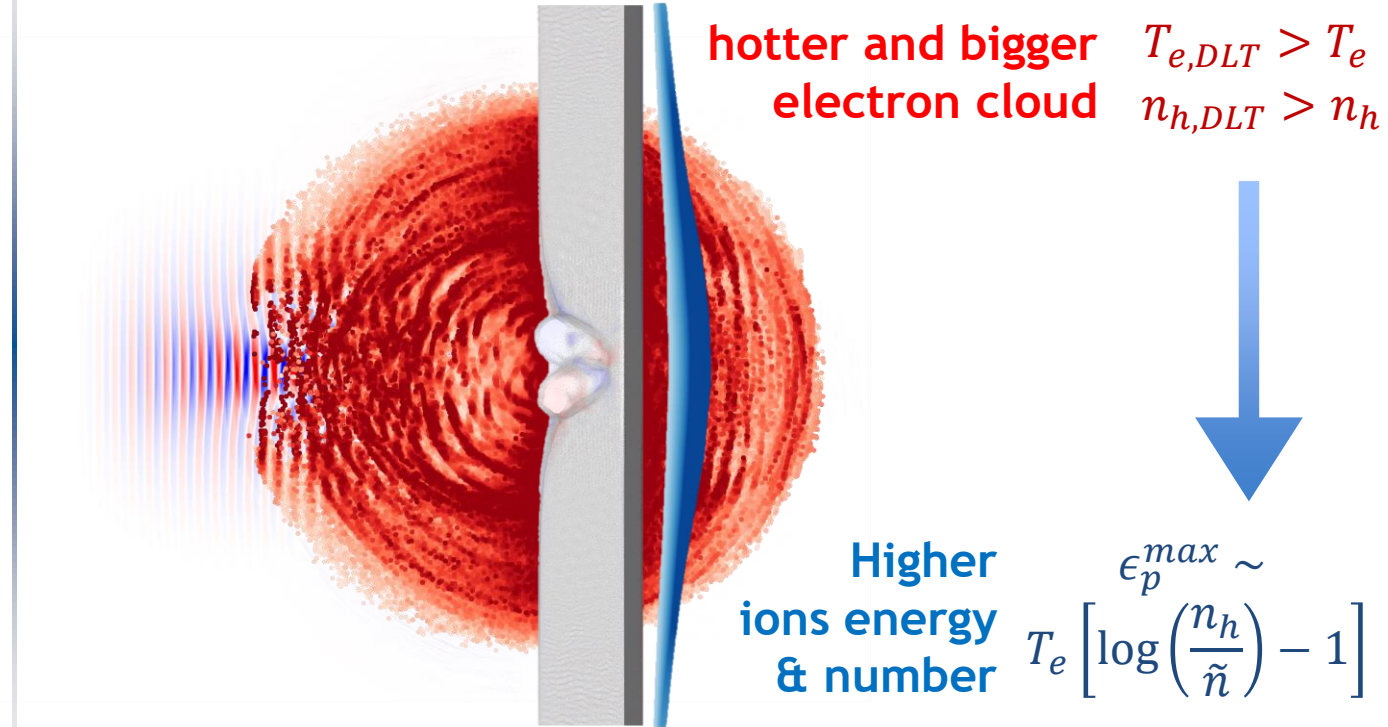


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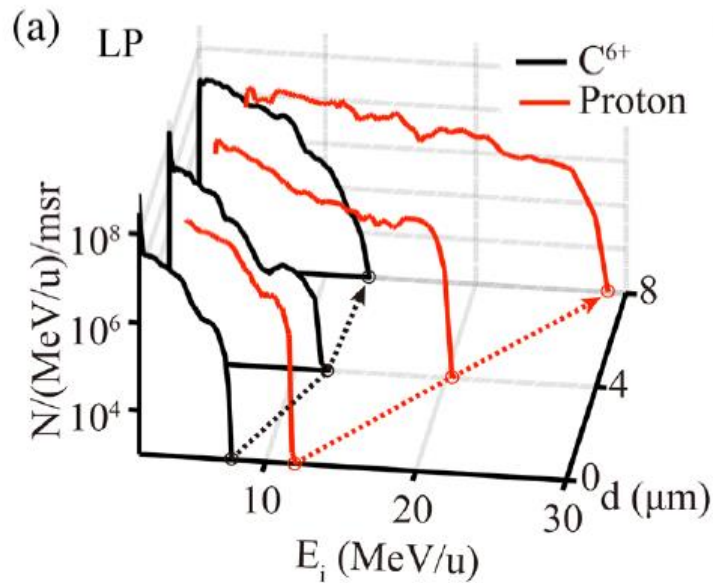
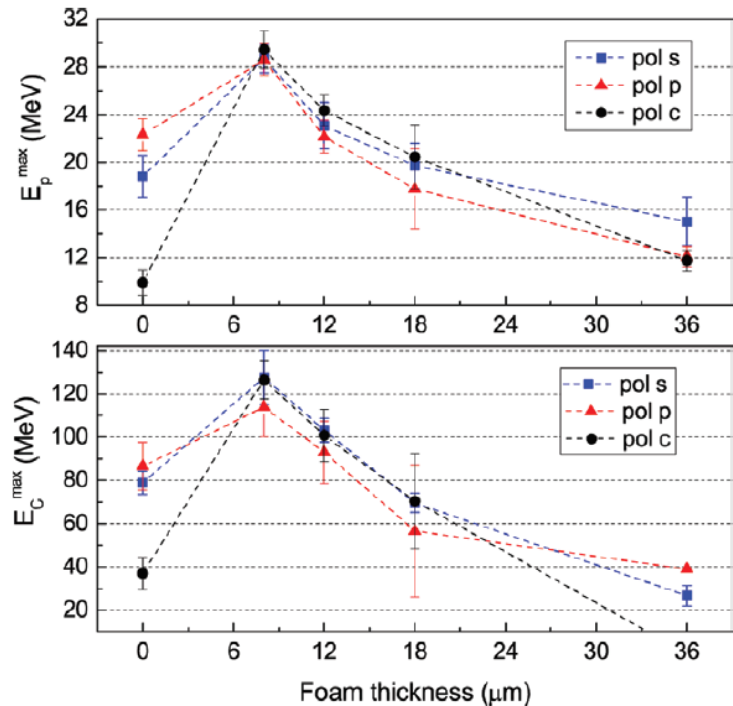
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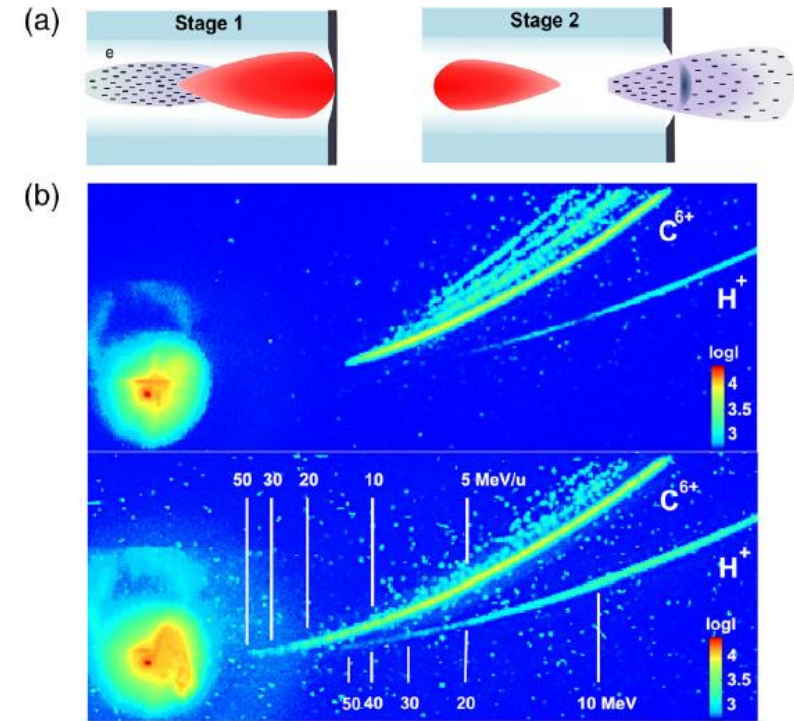
Advanced acceleration via near-critical DLT

Experimental evidences

Tested by independent groups:



Bin, J. H., et al. *Physical review letters* 120.7 (2018): 074801.



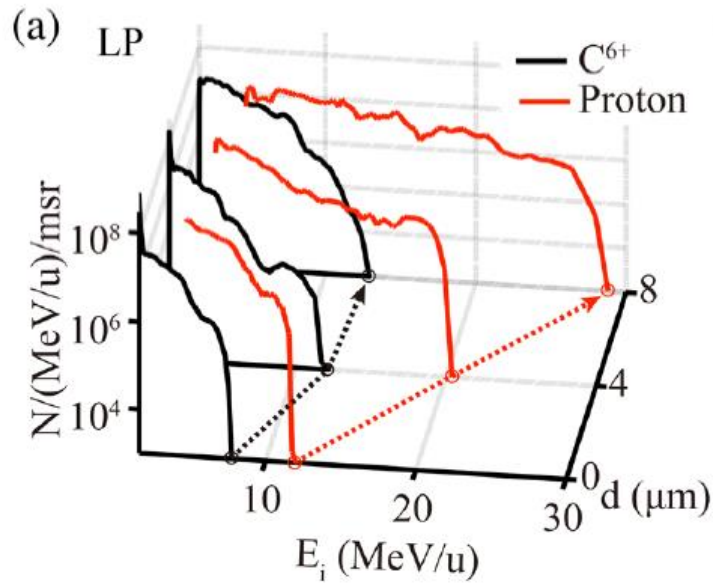
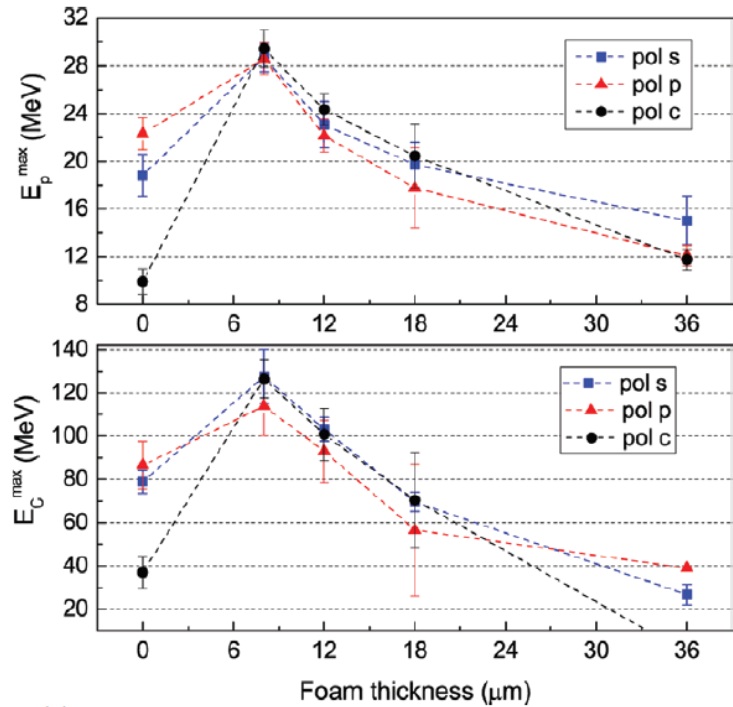
Ma, W. J., et al *Physical review letters* 122.1 (2019): 014803.

Prencipe, Irene, et al. *Plasma Physics and Controlled Fusion* 58.3 (2016): 034019.

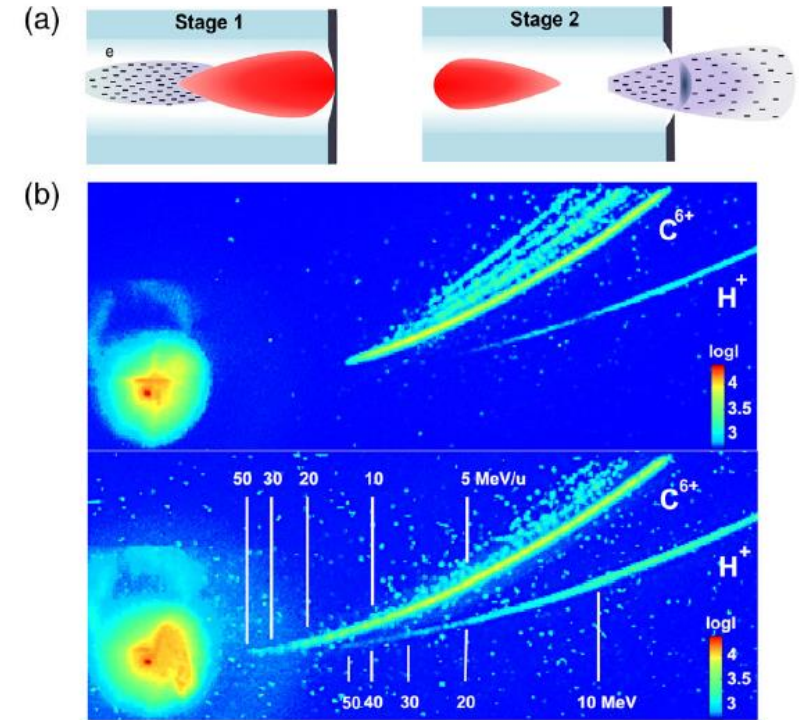
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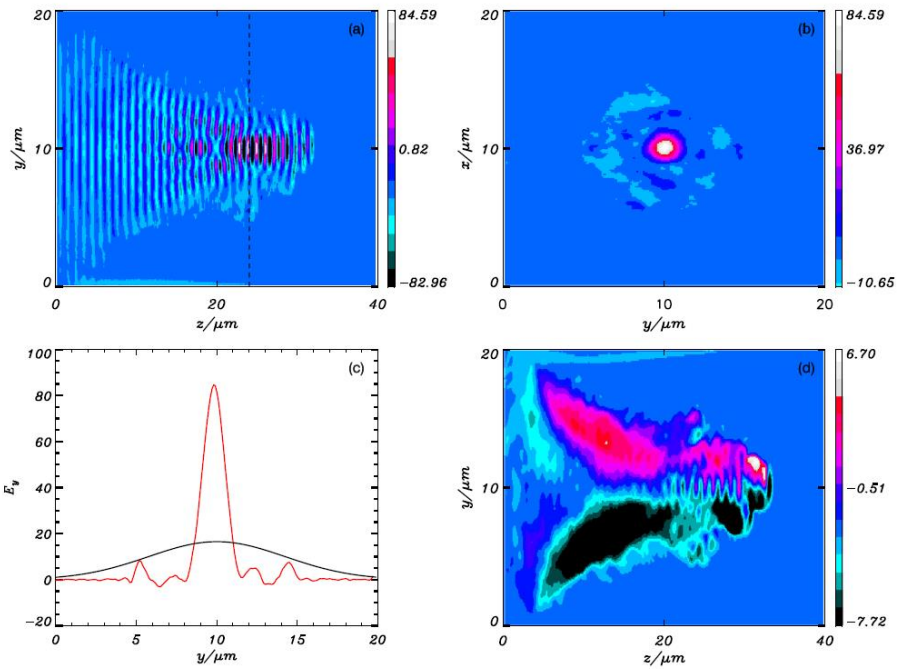
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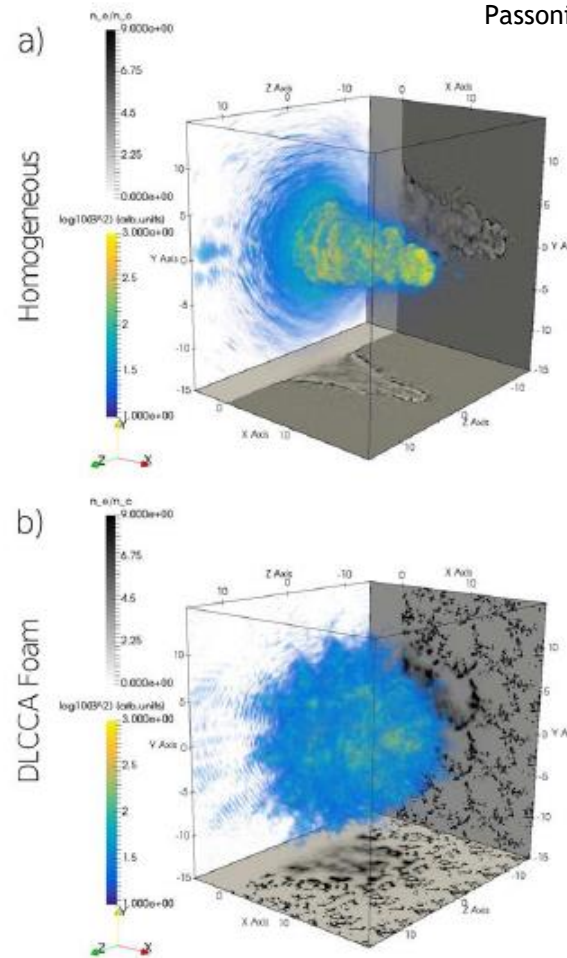
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Have we achieved the best performances?

PIC simulations help to understand the interaction physics

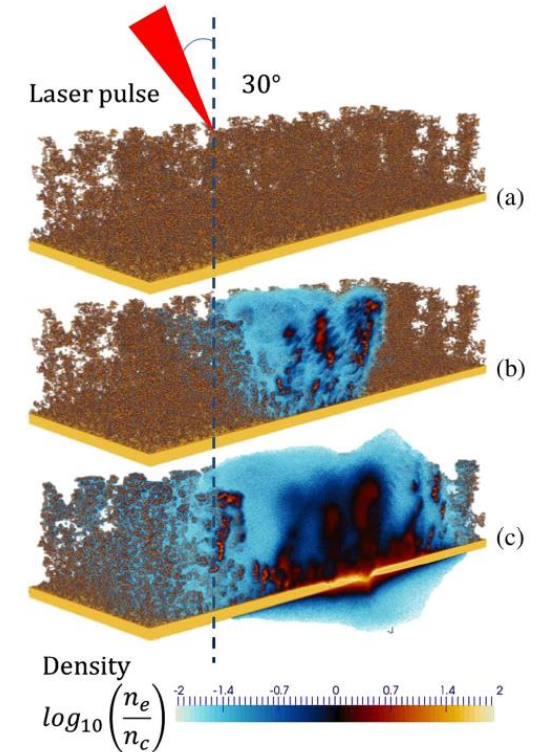


Wang, H. Y., et al. *Physical review letters* 107.26 (2011): 265002.

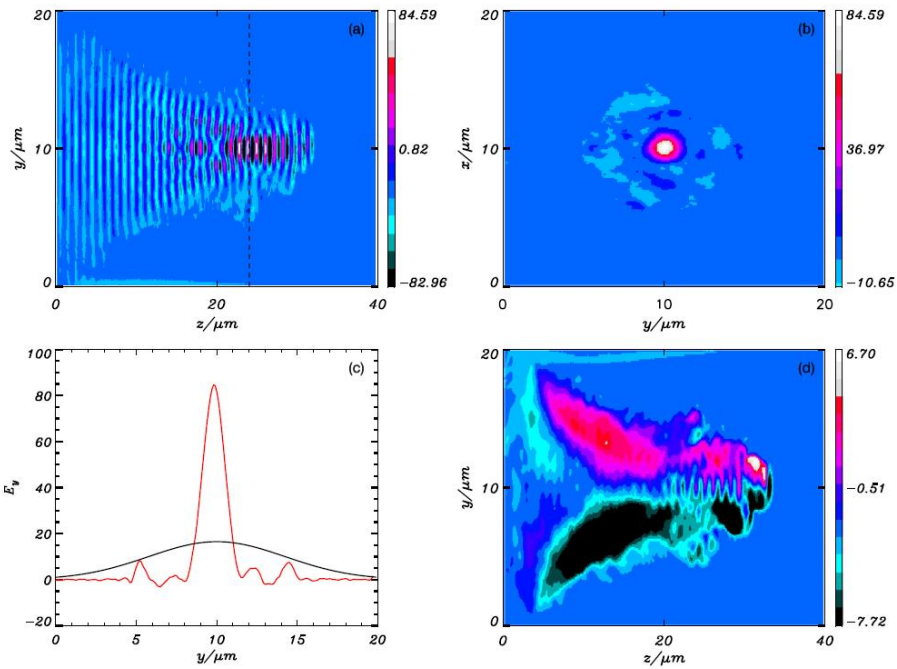


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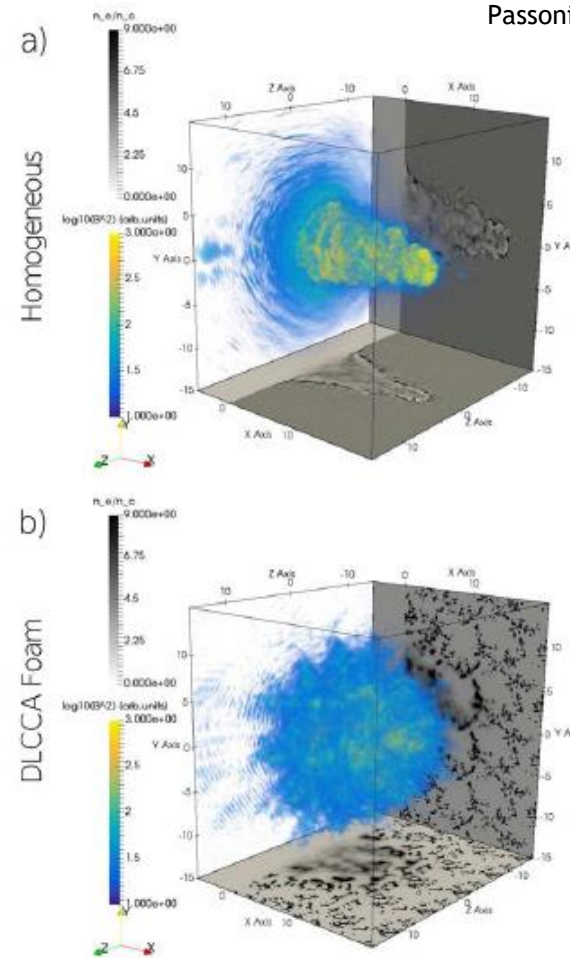
Fedeli, Luca, et al. *Scientific reports* 8.1 (2018): 3834.



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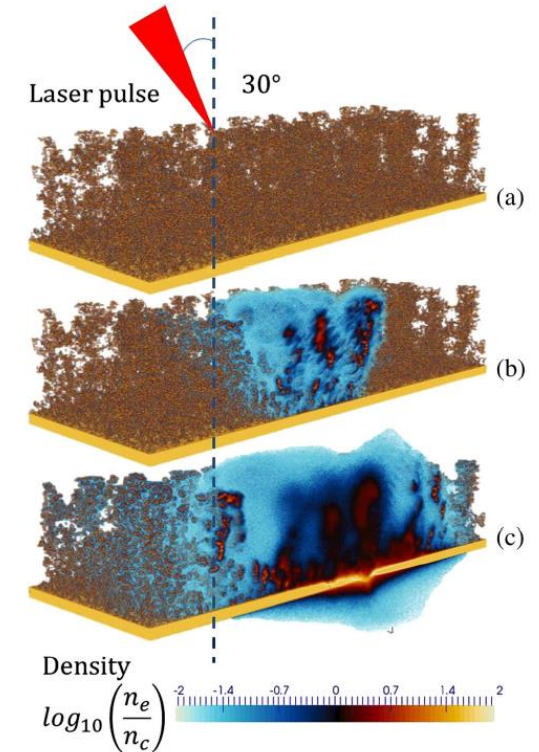


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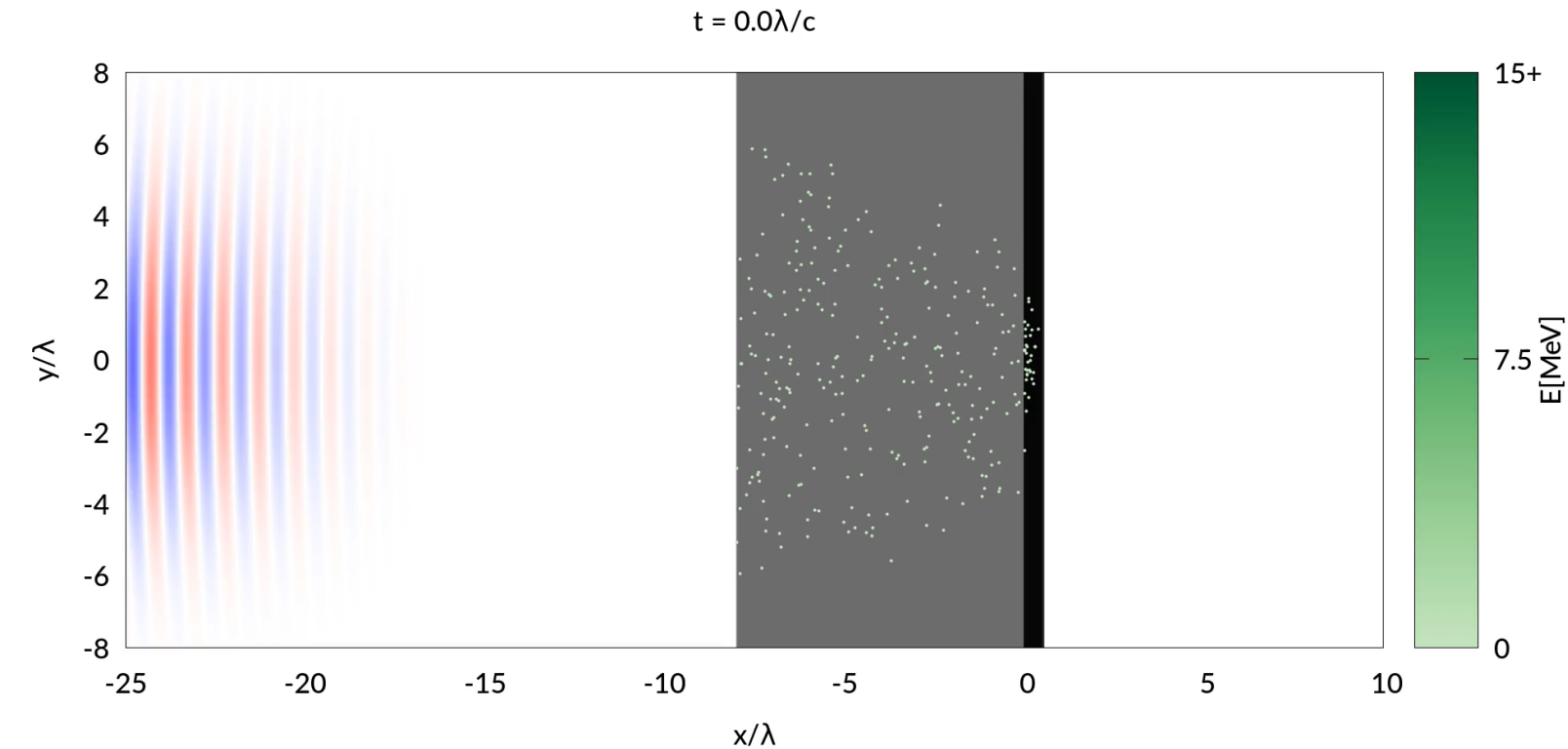
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Realistic 3D PIC simulations are computationally very expensive

From PIC simulations to a minimal model

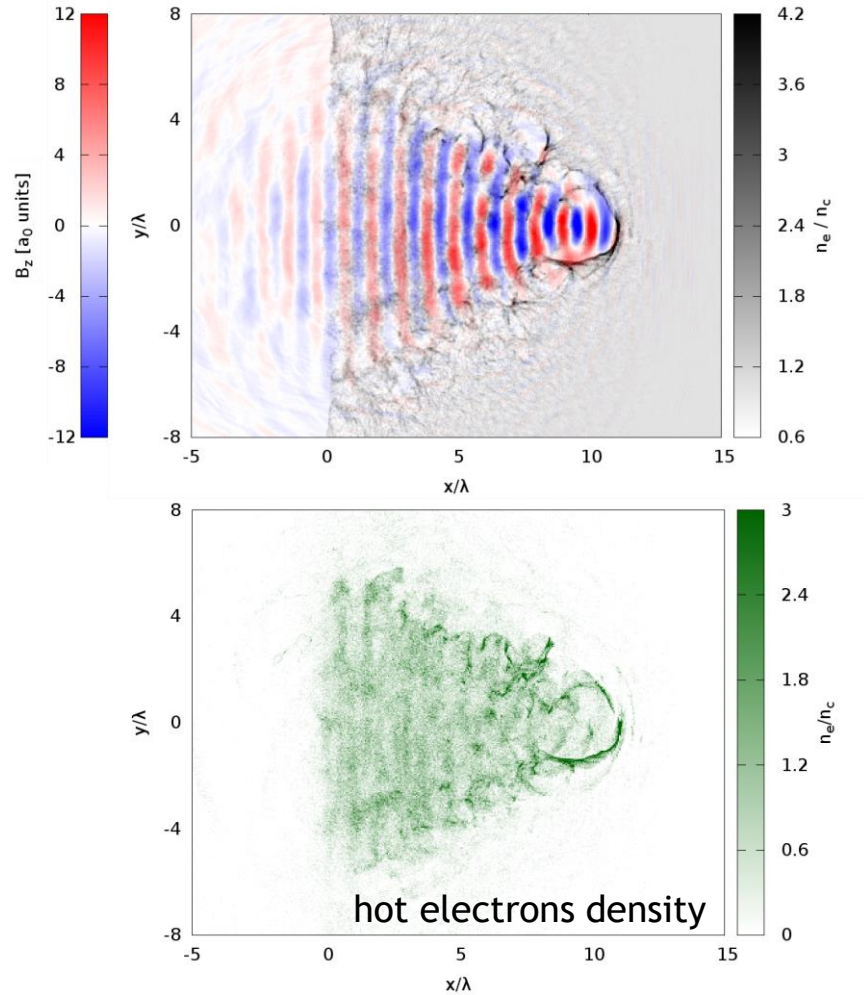


Main observed phenomena:

- Pulse drilling a channel
- Pulse self-focusing
- Hot electrons generation
- Self-generated magnetic fields

A modellistic approach is beneficial for the DLT optimization

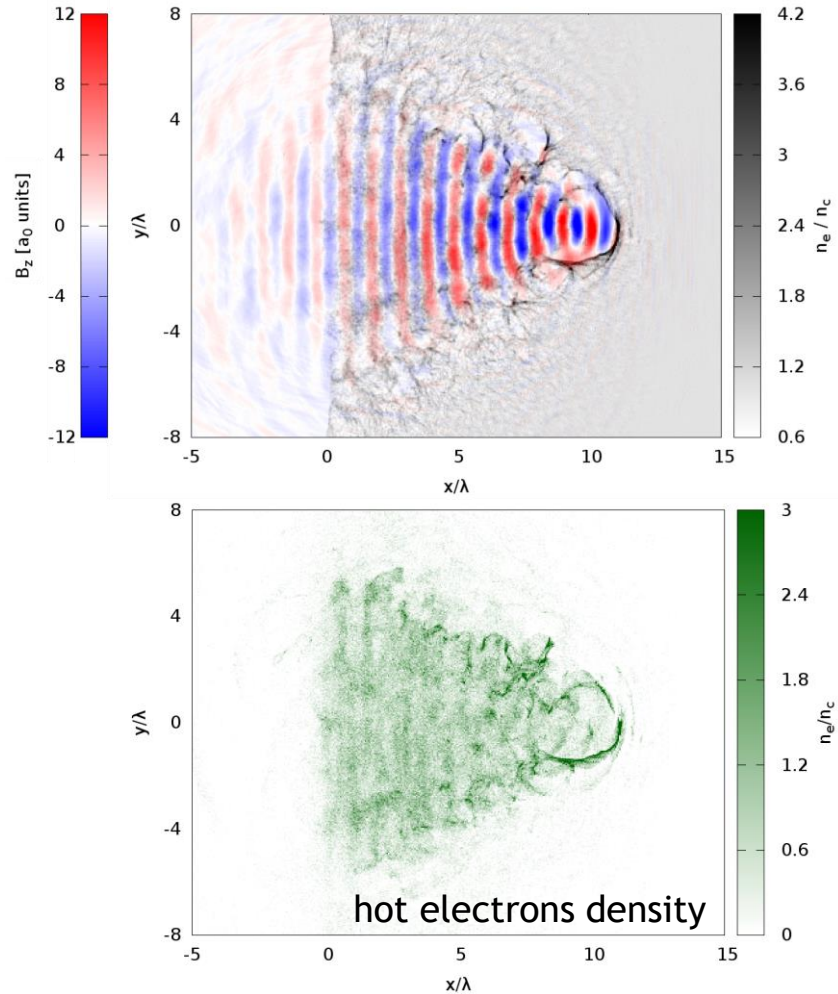
A minimal model is proposed



METHODS:

1. Theoretical model with free parameters
2. Parameters estimations with 2D/3D PIC simulations

A minimal model is proposed



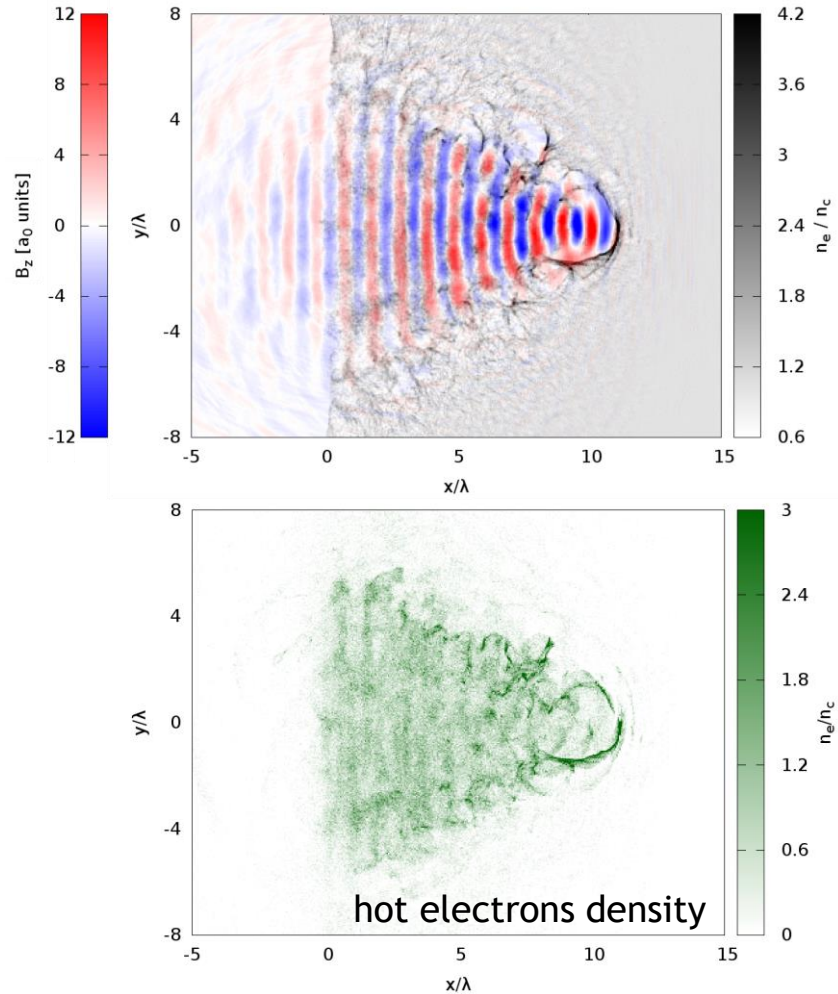
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MODEL STEPS:

1. Laser self-focusing
2. Laser energy loss and amplification
3. Hot electrons heating
4. Ions acceleration

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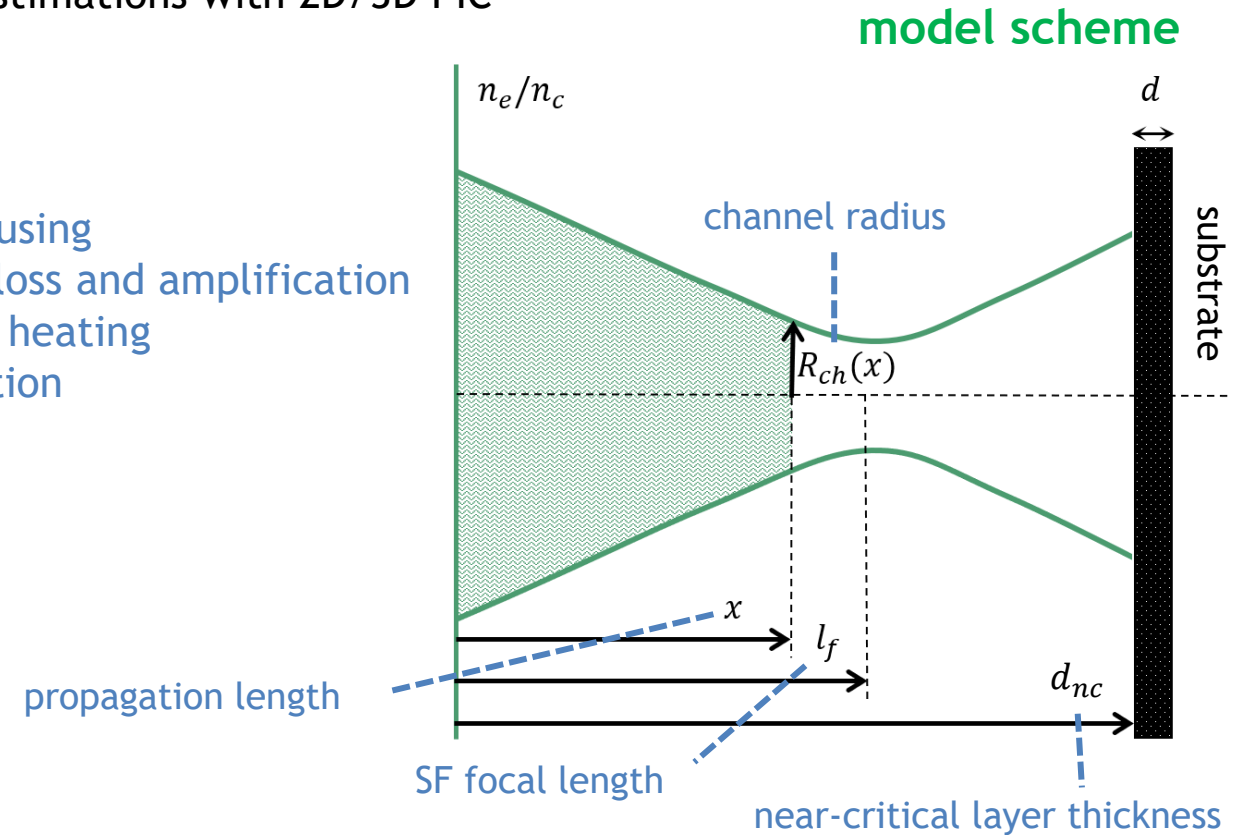


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MODEL STEPS:

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1st step: laser propagation into a near-critical plasma

Pulse waist focuses with
a thin-lens law:

$$w(x) = w_m \sqrt{1 + \left(\frac{x - l_f}{x_R}\right)^2}$$

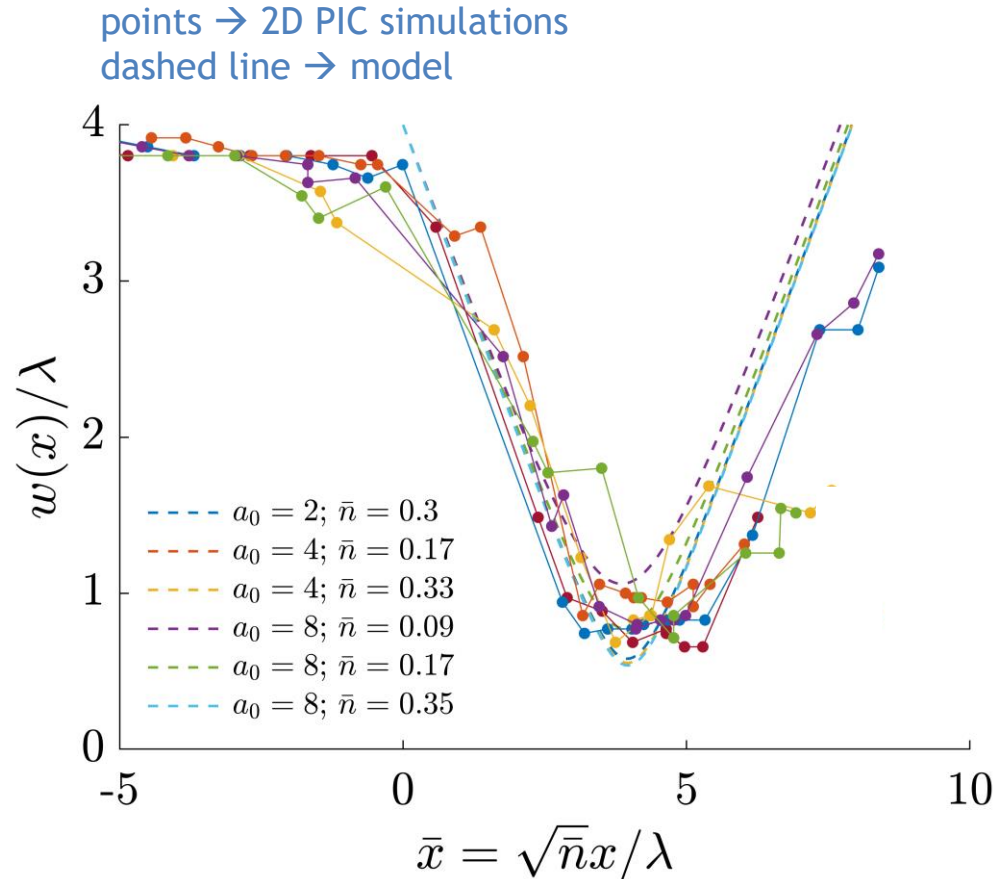
minimum waist

$$w_m = \frac{\lambda}{\pi} \frac{1}{\sqrt{\bar{n}}}$$

Wang, H. Y., et al. *Physical review letters* 107.26 (2011): 265002.

relativistic
transparency factor

$$\bar{n} = \frac{n_e}{\gamma n_c}$$



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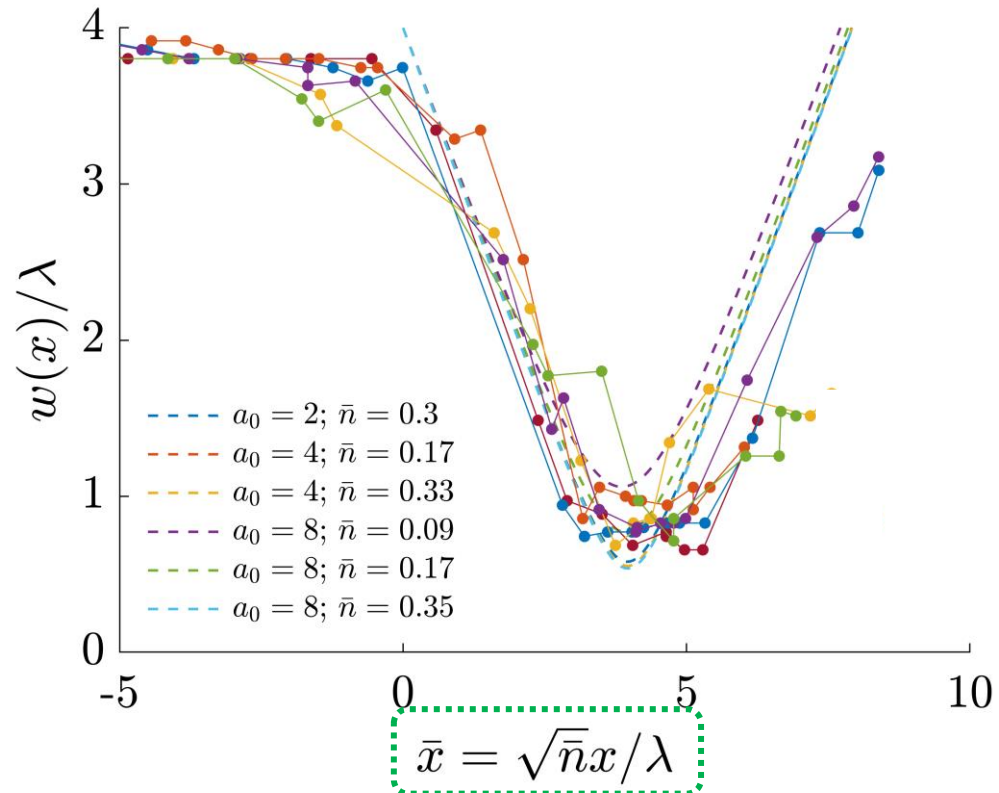
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points → 2D PIC simulations
dashed line → model



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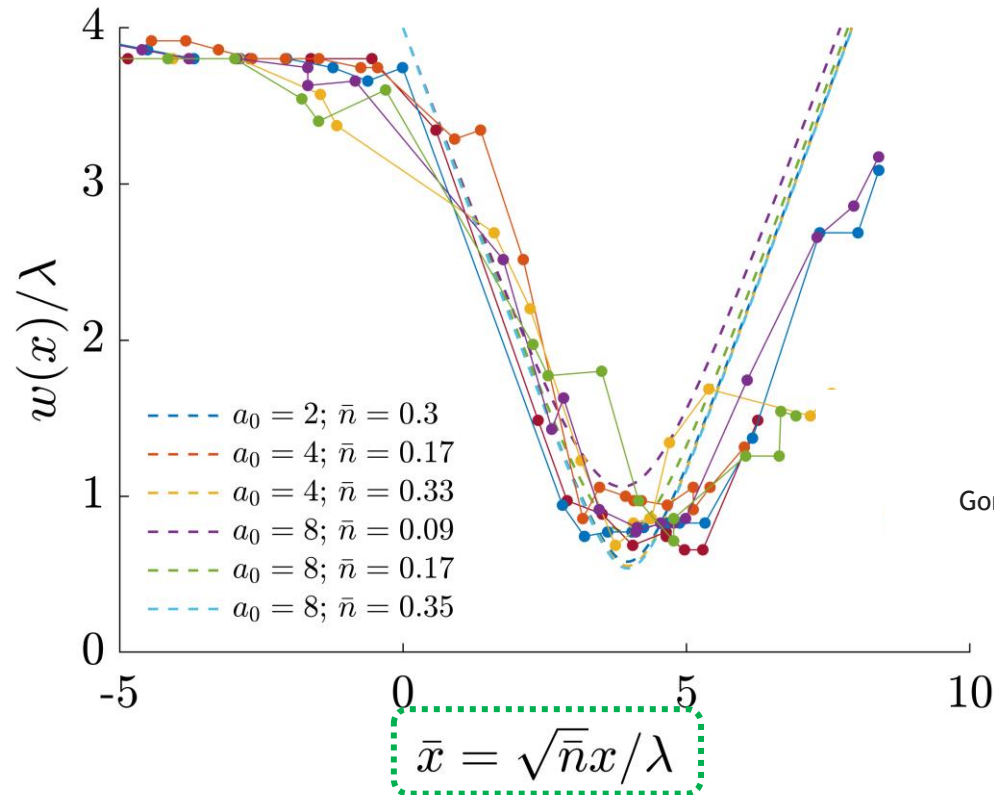
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thin-lens law can be written as:

$$w(\bar{x}) \sim \lambda \sqrt{\frac{1}{\pi^2 \bar{n}} + \left(\bar{x} - \frac{w_0}{\lambda}\right)^2}$$

↓
 $\bar{x} = \sqrt{\bar{n}} x / \lambda$

The propagation is self-similar!

Gordienko, S., and A. Pukhov. *Physics of Plasmas* 12.4 (2005): 043109.

2nd step: pulse energy loss and amplification

pulse energy loss equation:

$$d\varepsilon_p = -T_{nc}(x)n_e 2R_{ch}(x)dx$$

← coupled with →

amplification equation:

$$\bar{a}(x) = \frac{a(x)}{a_0} = \sqrt{\frac{\bar{\varepsilon}_p(x)}{w(x)/w_0}}$$

2nd step: pulse energy loss and amplification

pulse energy loss equation:

$$d\varepsilon_p = -T_{nc}(x)n_e 2R_{ch}(x)dx$$

electron temperature $T_{nc}(x) = C_{nc}[\gamma(x) - 1]m_e c^2$ **corrected ponderomotive scaling**
Cialfi, Lorenzo, Luca Fedeli, and Matteo Passoni. *Physical Review E* 94.5 (2016): 053201.

channel radius $R_{ch}(x) = r_c w(x)$ **proportional to pulse waist**

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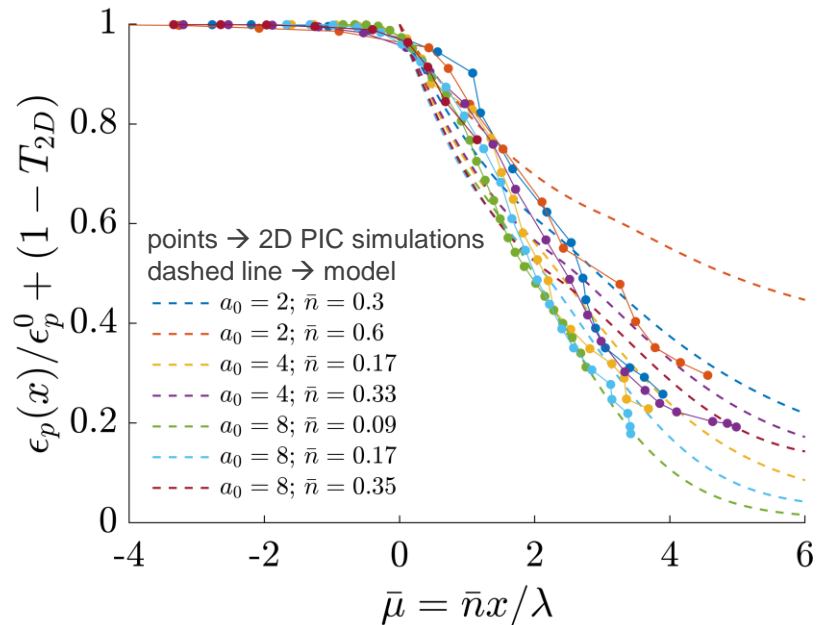
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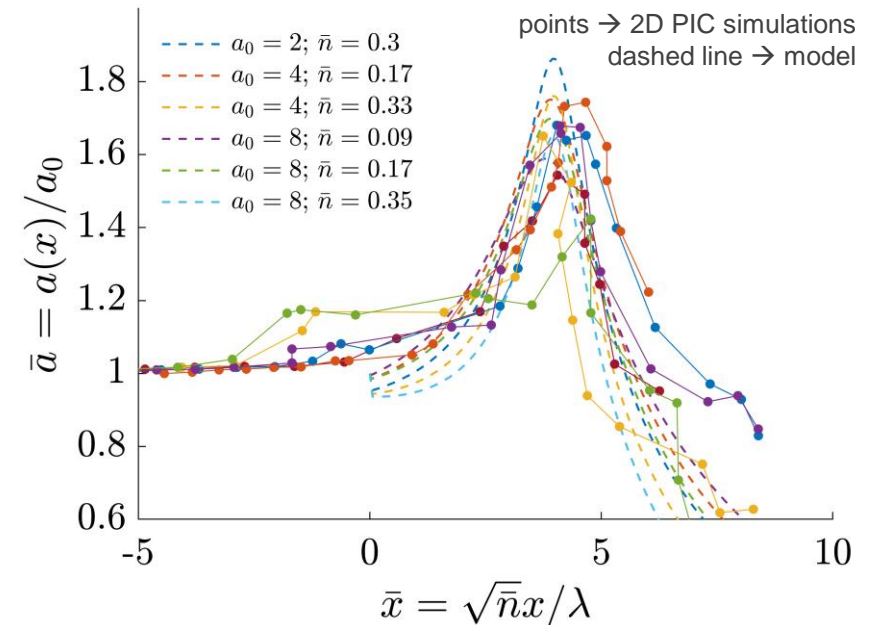
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Two free parameters: $C_{nc} = 1.7; r_c = 2.0$

3rd step: hot electrons heating

The energy lost by the pulse is given to the hot electrons:

$$E_{nc}(x) = \frac{(1 - \overline{\varepsilon}_p(x)) \varepsilon_{p0}}{N_{nc}(x)}$$

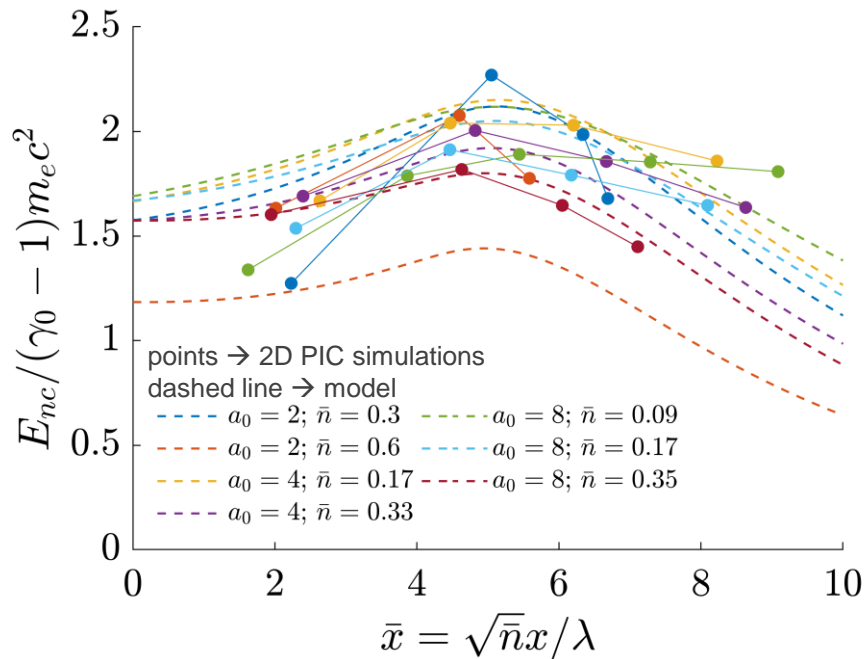
→ Total number of electrons
in the channel

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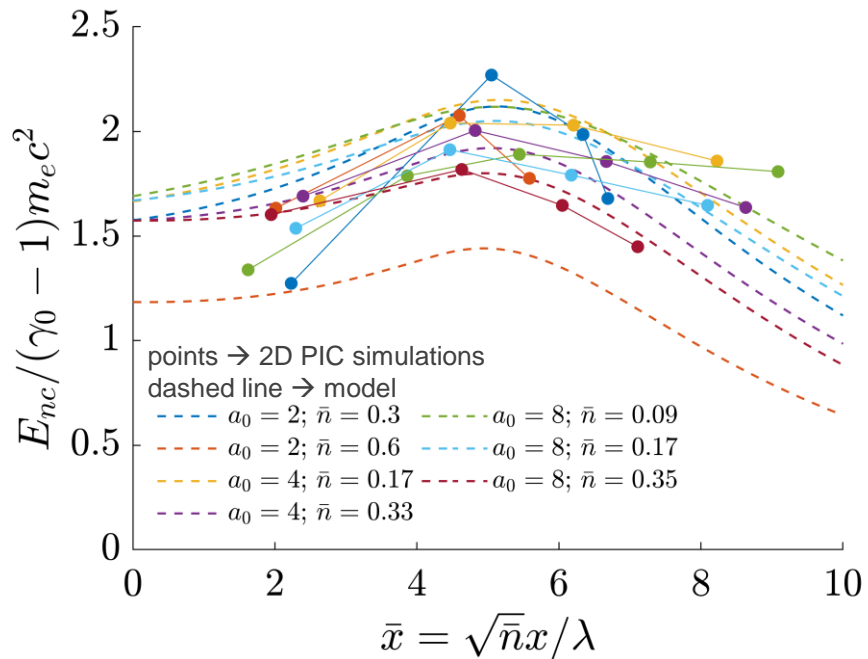


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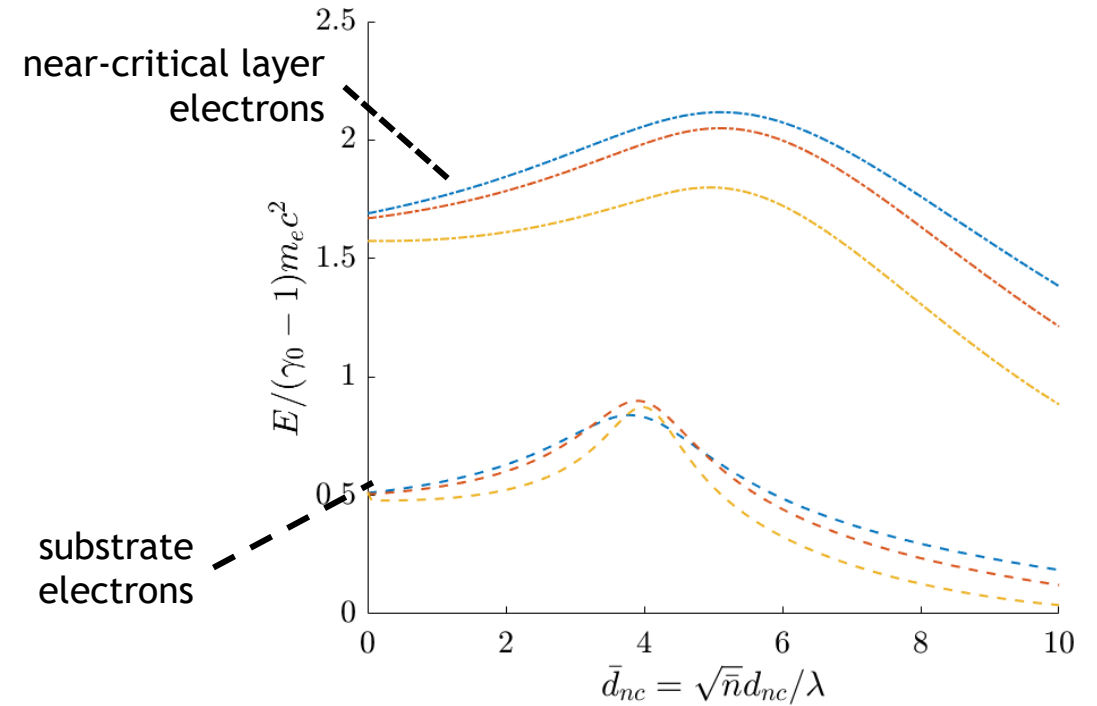
$$E_{nc}(x) = \frac{(1 - \overline{\varepsilon}_p(x)) \varepsilon_{p0}}{N_{nc}(x)}$$

$N_{nc}(x)$ → Total number of electrons in the channel



Also electrons from the substate are considered:

$$E_s(d_{nc}) = C_s [\gamma(d_{nc}) - 1] m_e c^2$$



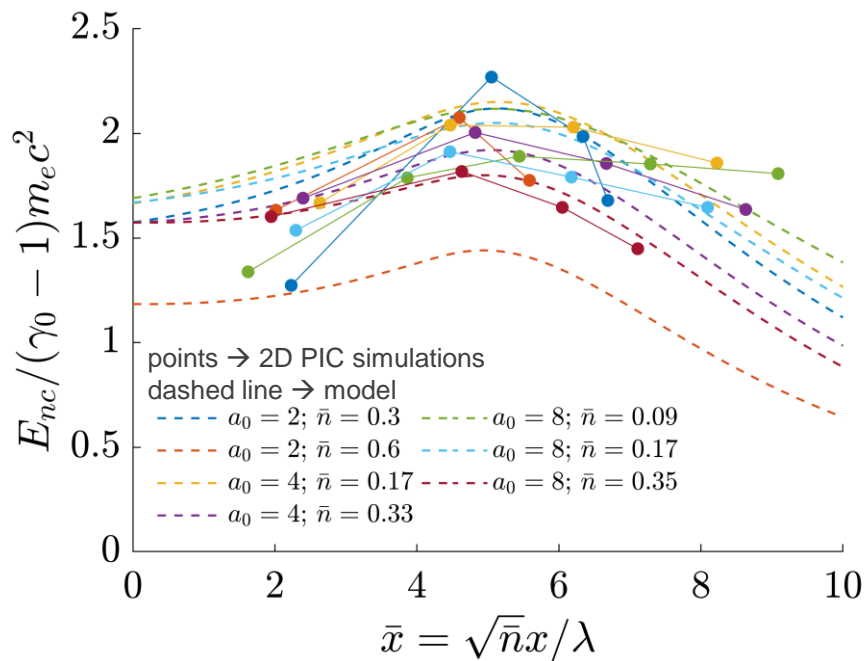
Additional free parameter: $C_s = 0.24$

3rd step: hot electrons heating

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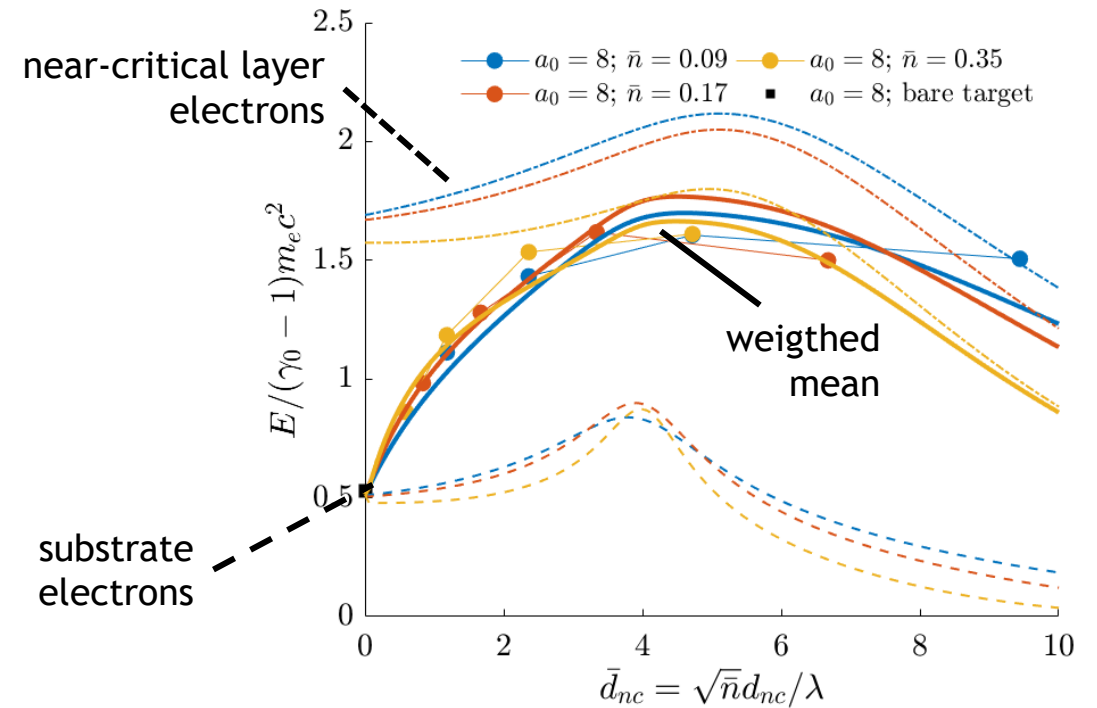
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4th step: proton maximum energy estimation

Quasi-static model:

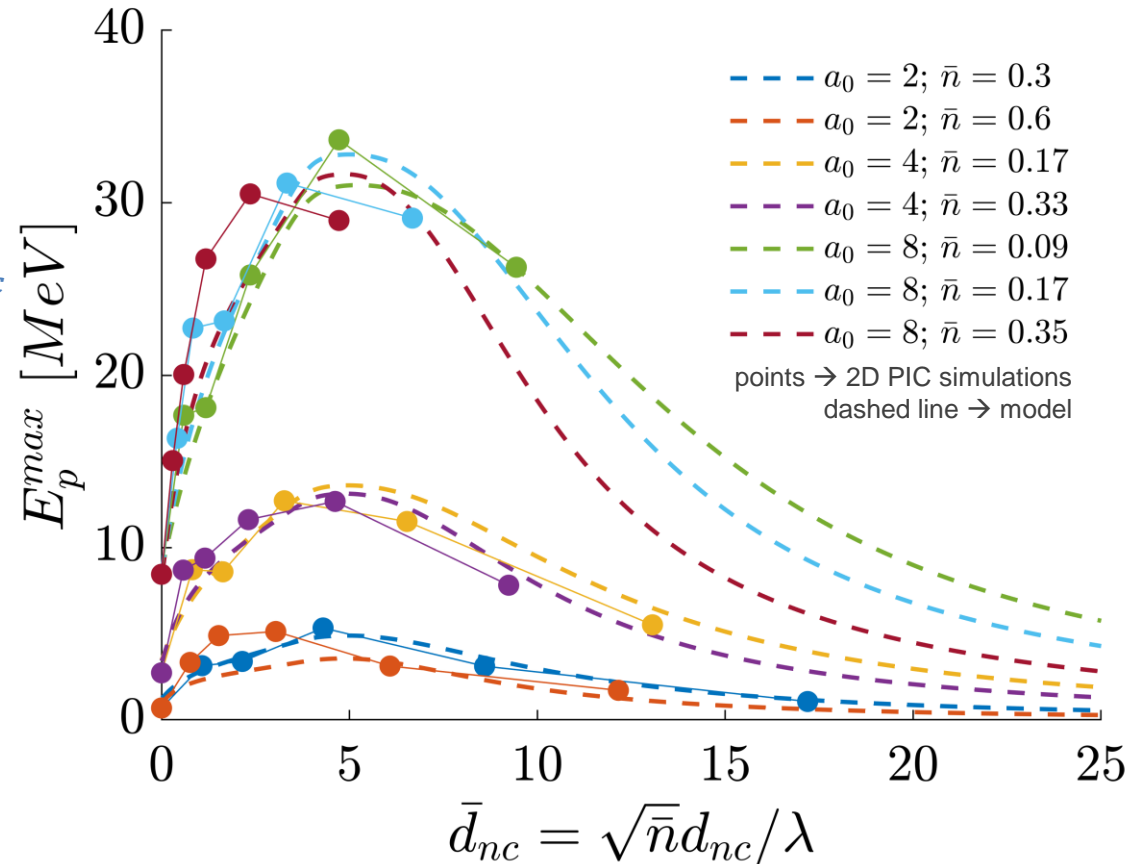
$$\epsilon_p^{max} = E_{DLT} \left[\log \left(\frac{n_{h,DLT}}{\tilde{n}} \right) - 1 \right]$$

4th step: proton maximum energy estimation

Quasi-static model:

$$\epsilon_p^{max} = E_{DLT} \left[\log \left(\frac{n_{h,DLT}}{\tilde{n}} \right) - 1 \right]$$

last free parameter: $\tilde{n} = 1.3 \cdot 10^{-3} n_c$

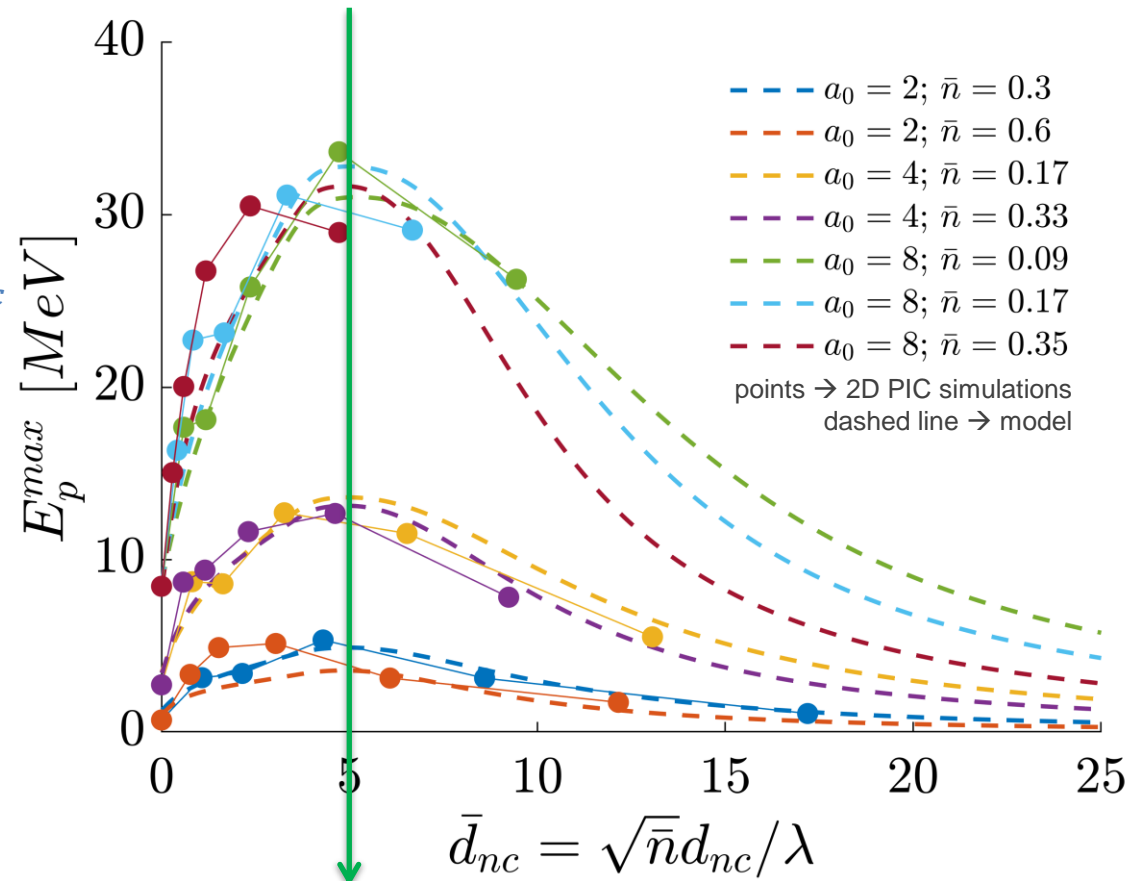


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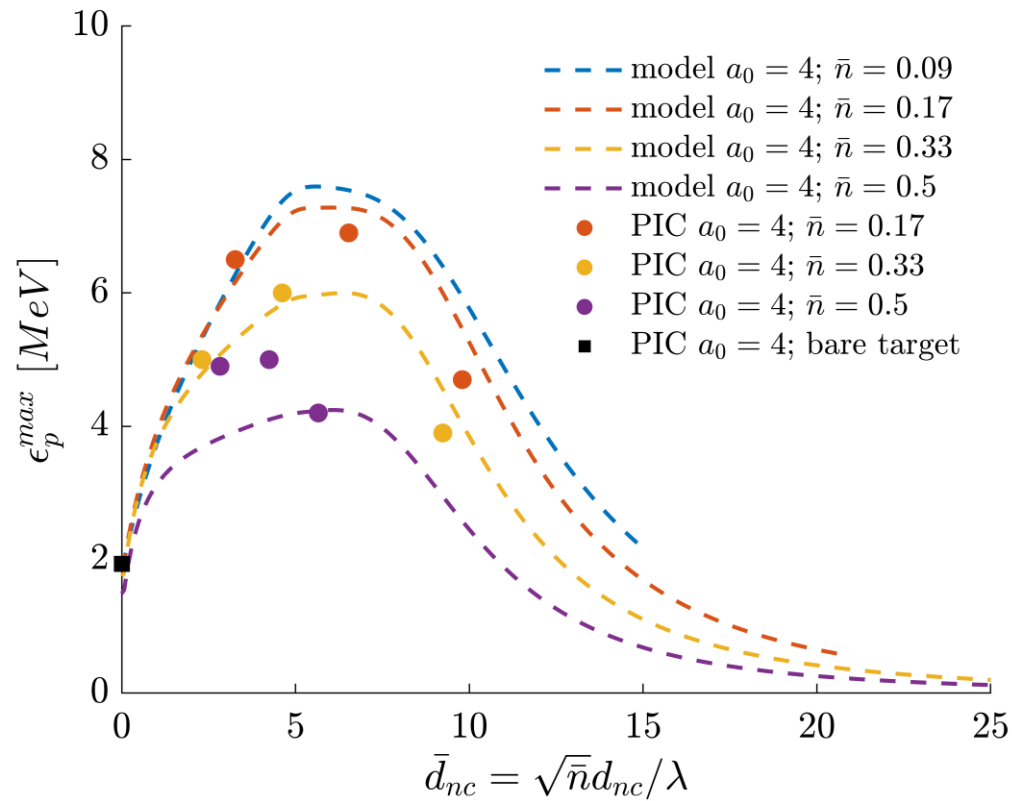


$$d_{nc}^{opt} \sim w_0 / \sqrt{\bar{n}}$$

The same model is solved in 3D

3D parameters:

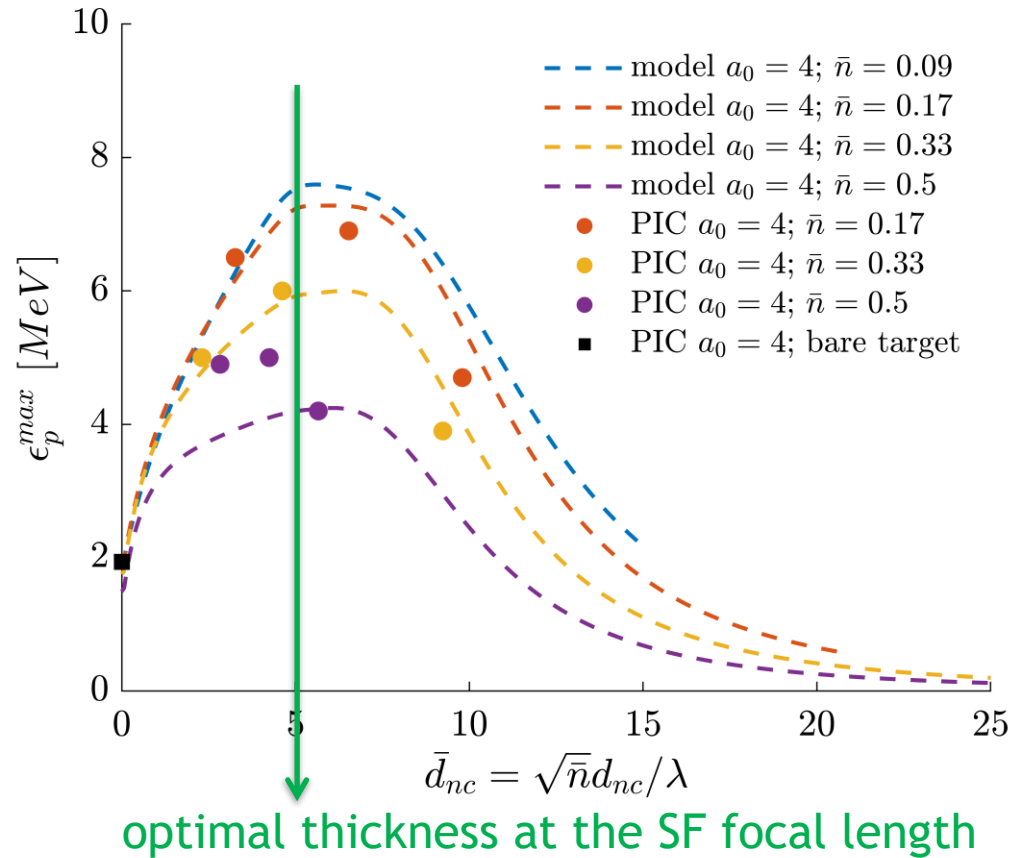
$$C_{nc} = 1.1$$
$$r_c = 2.1$$
$$C_s = 0.18$$
$$\tilde{n} = 5 \cdot 10^{-2} n_c$$



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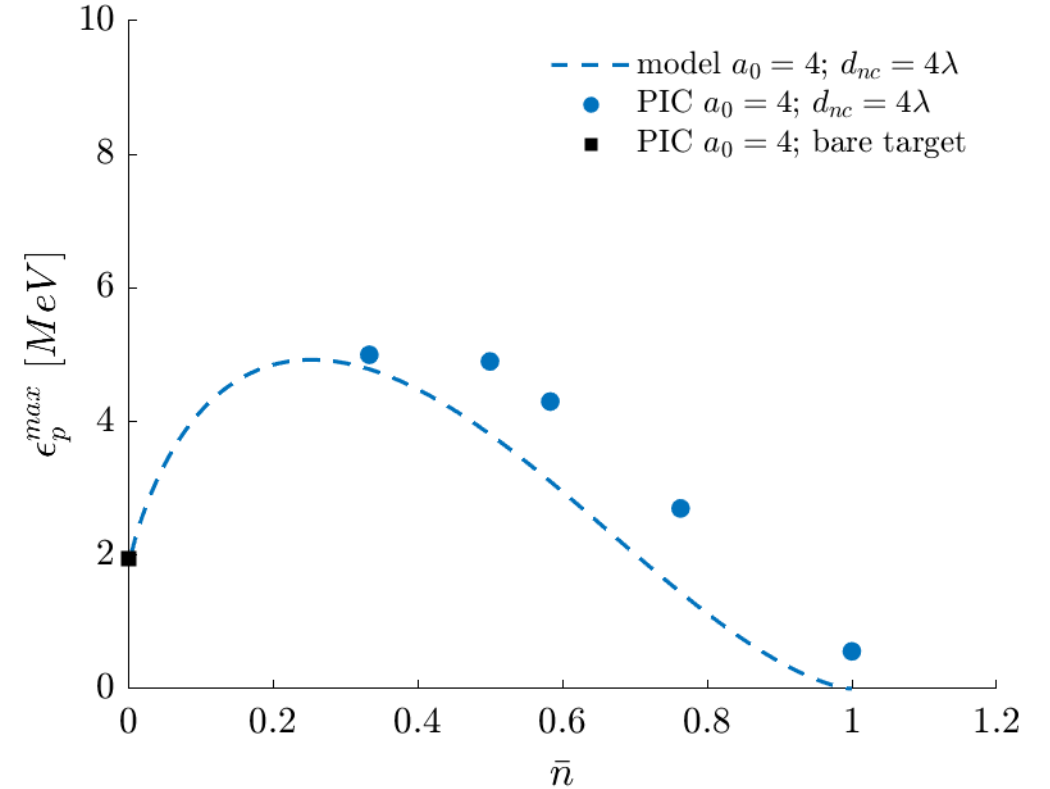
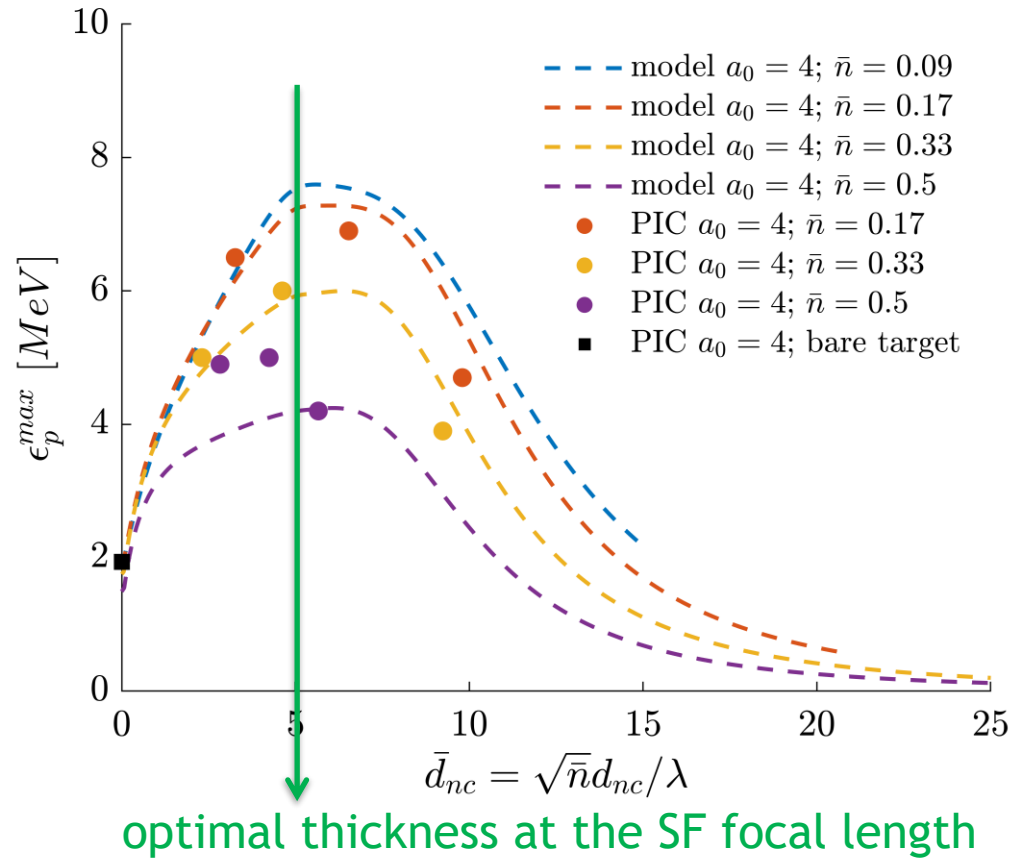
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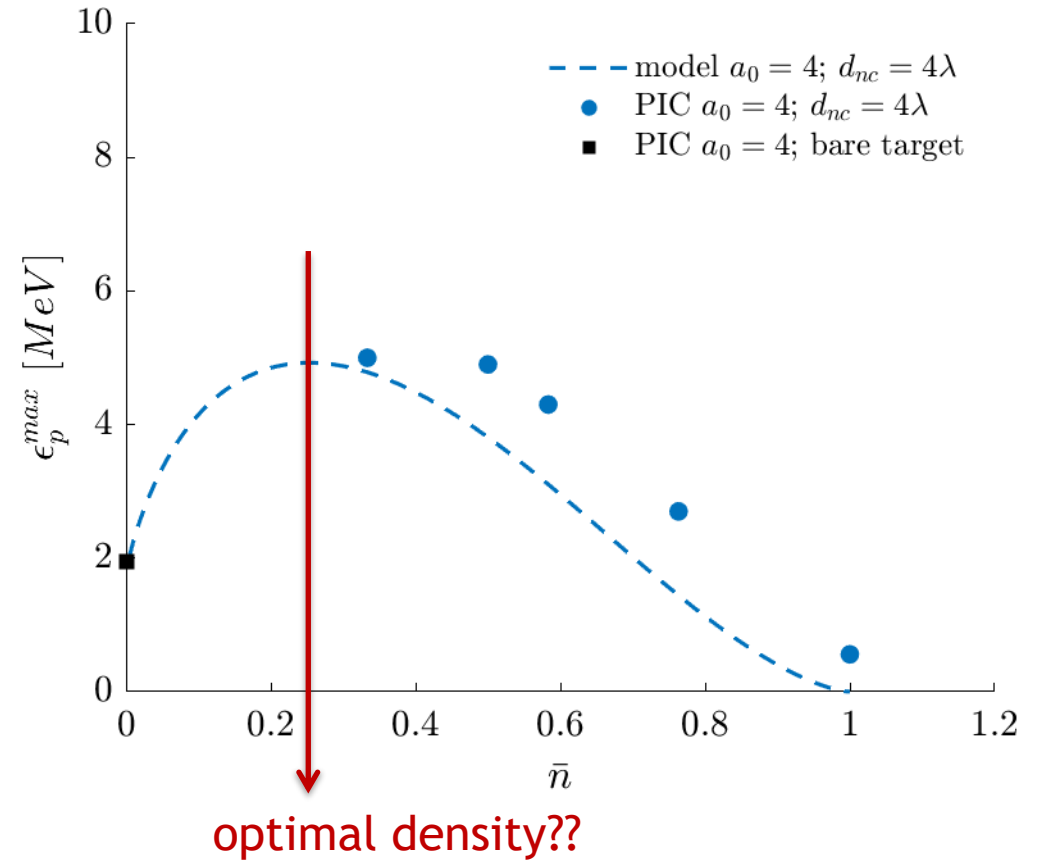
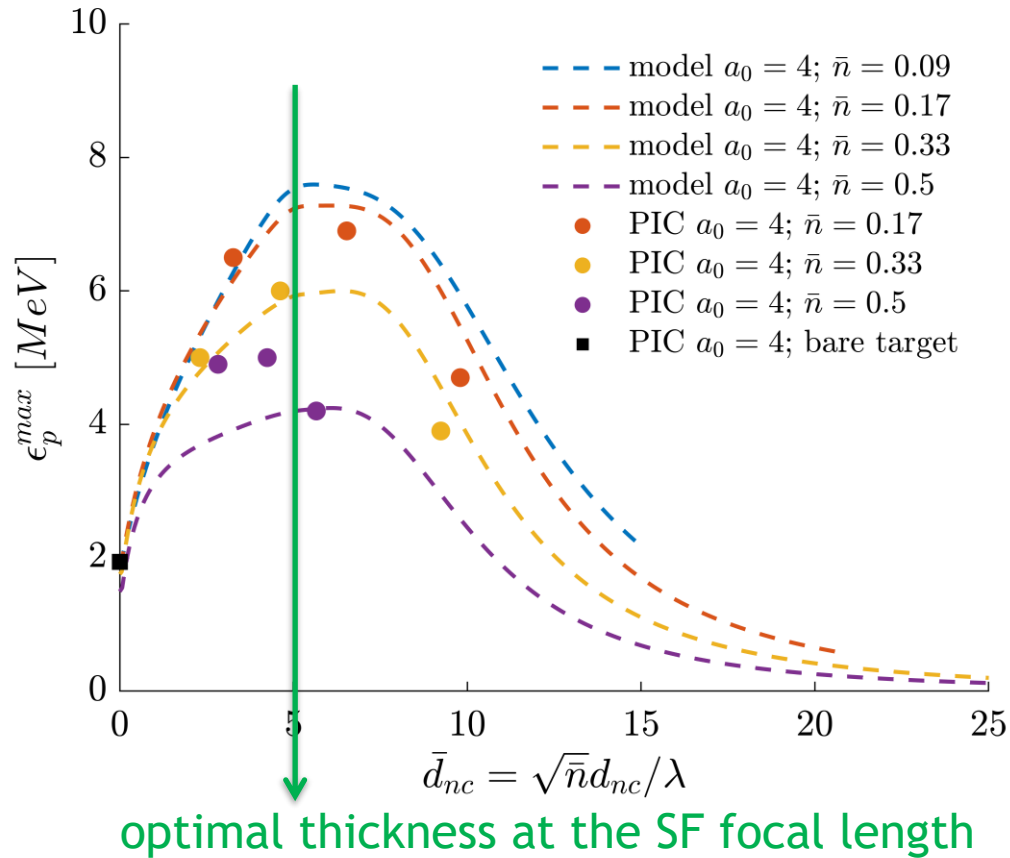
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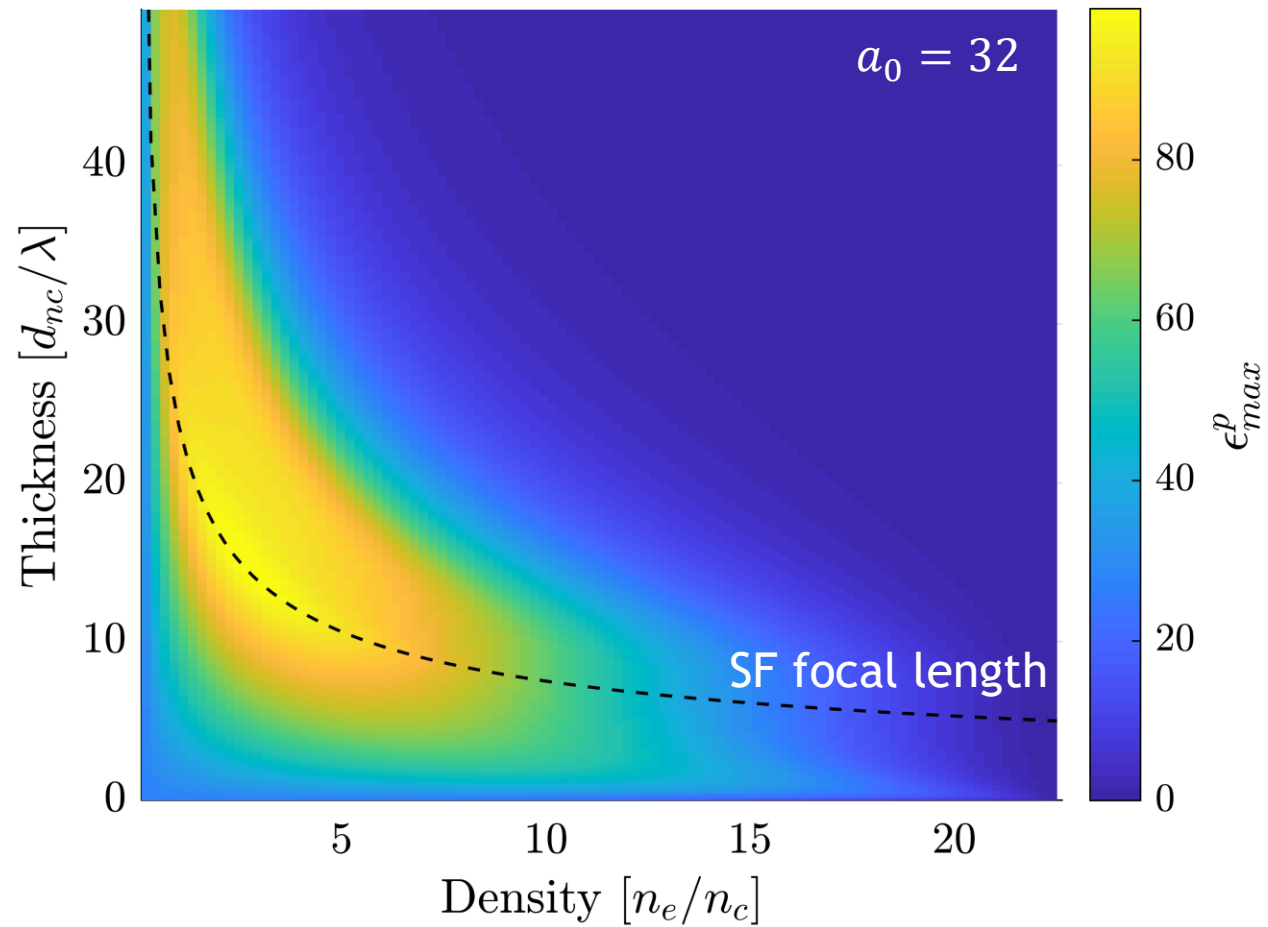
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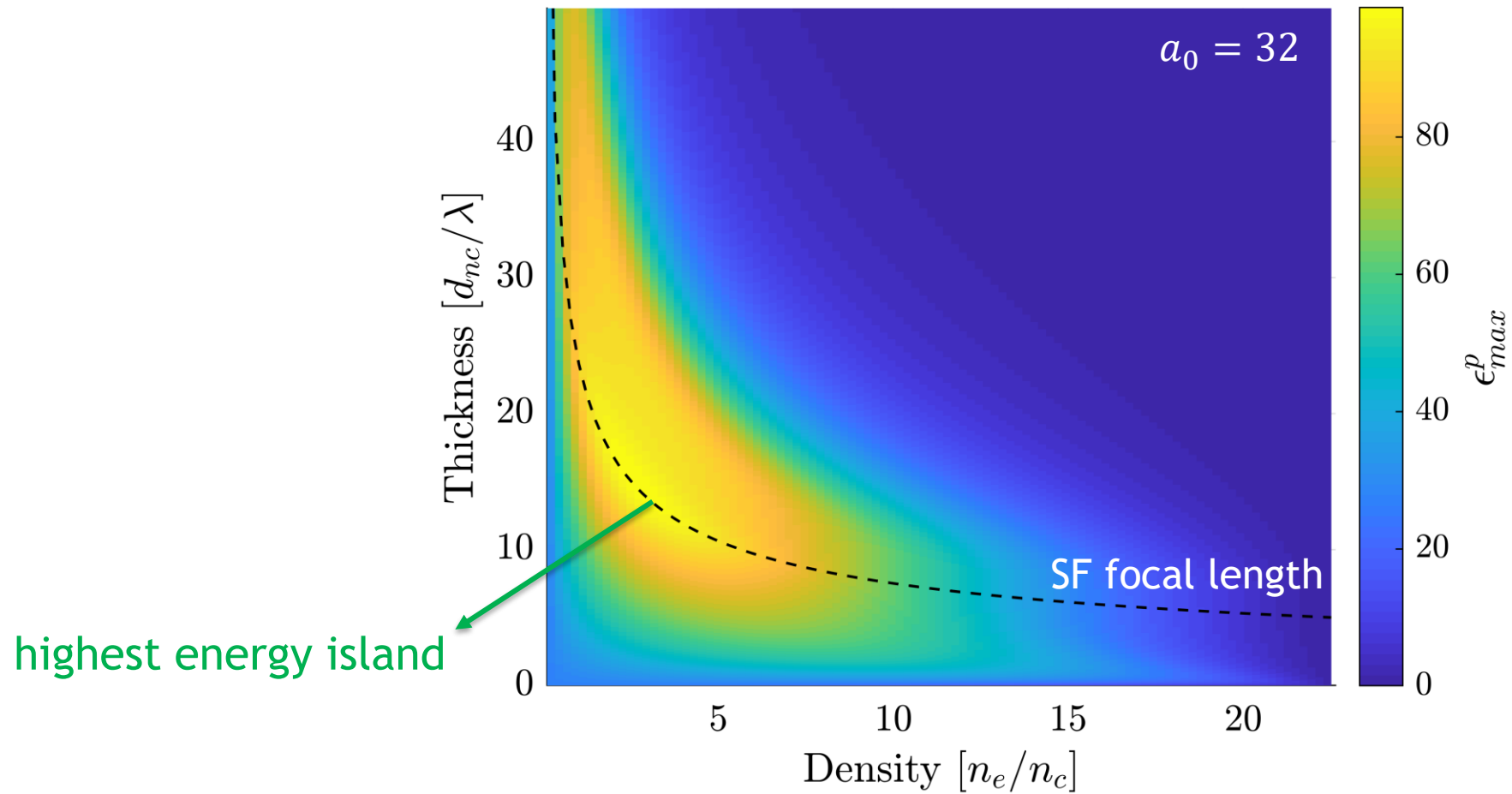
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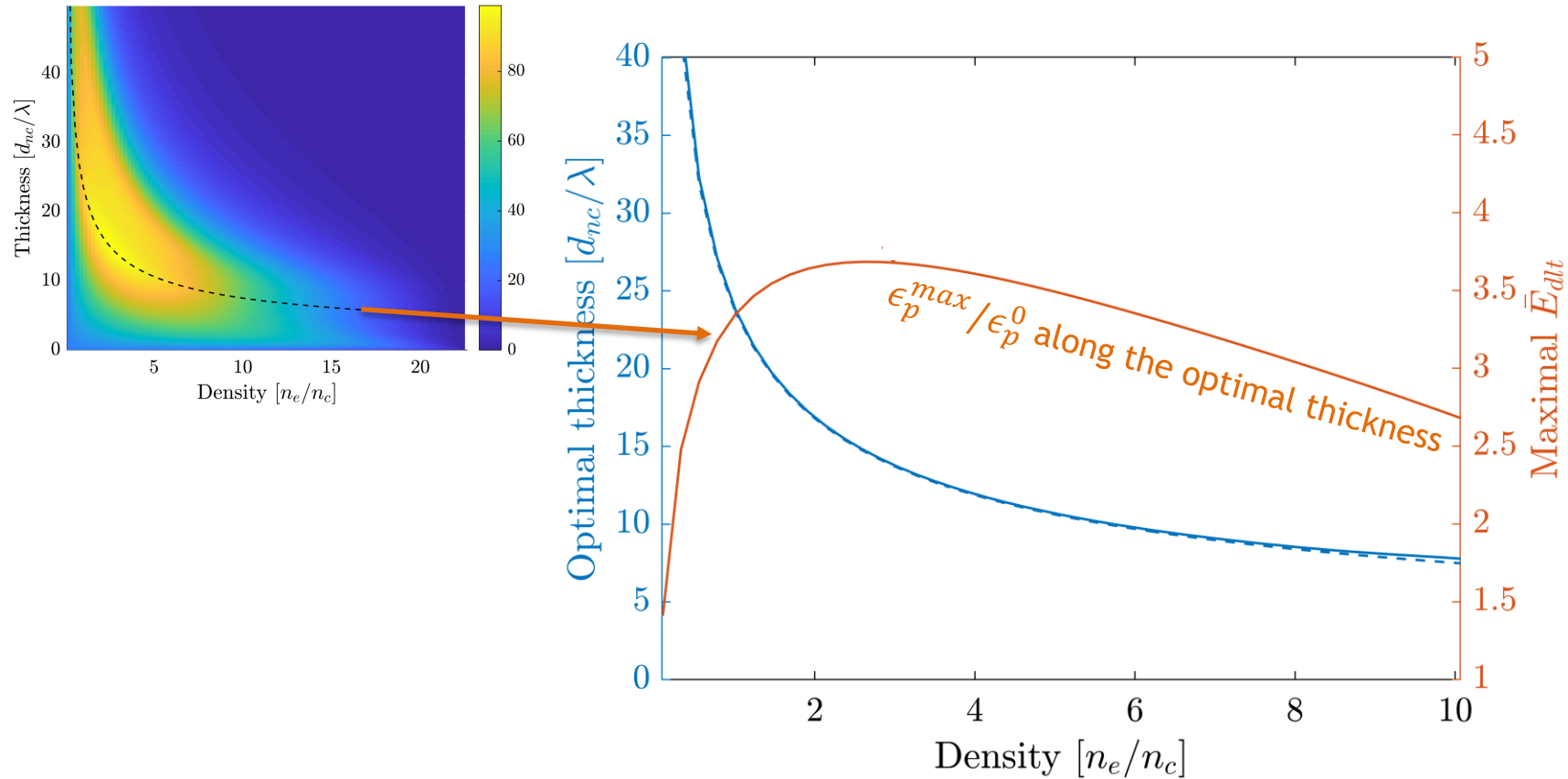
Maximum proton energy heat map



Maximum proton energy heat map

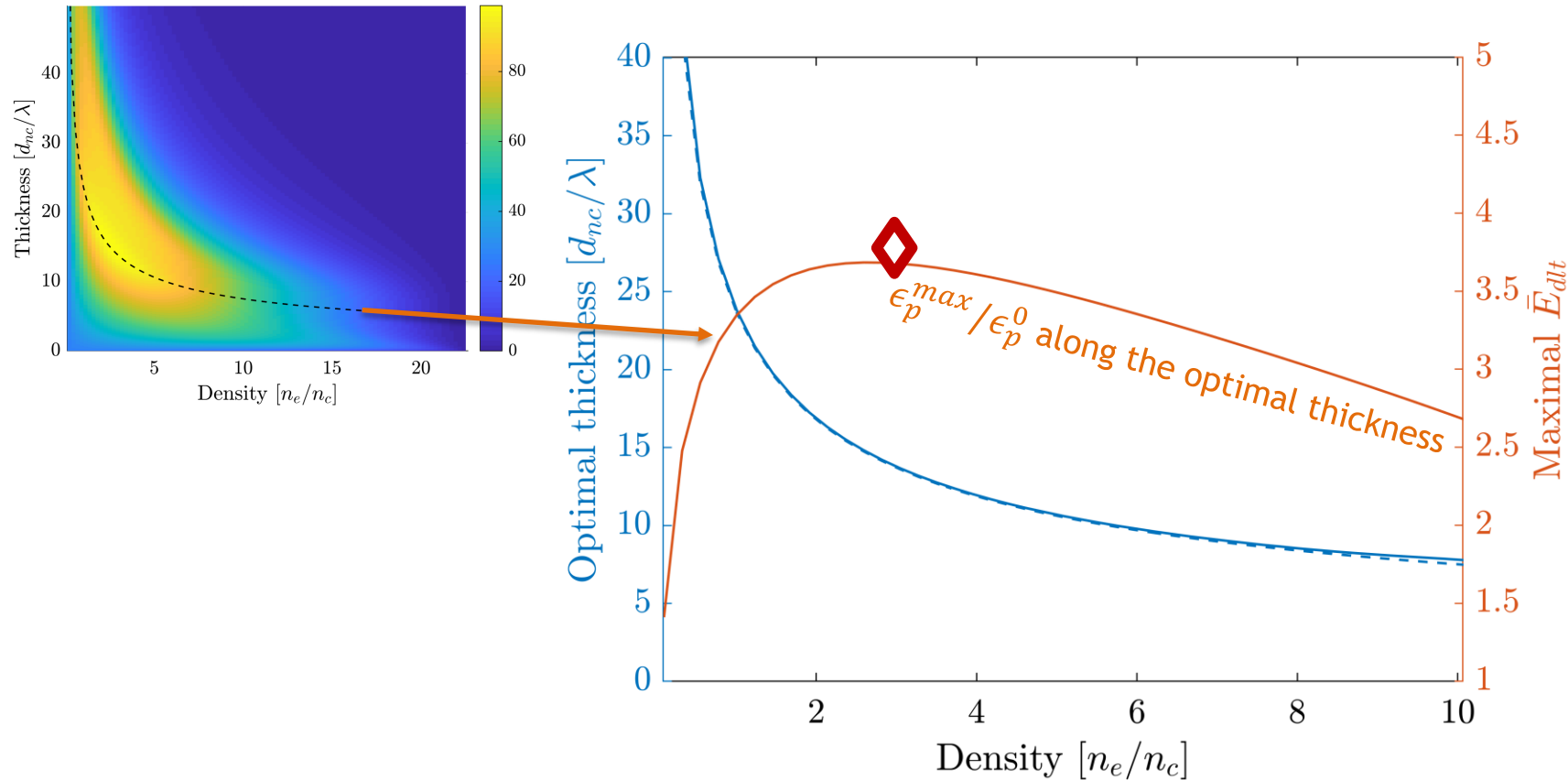


Analytical solution in the ultra-relativistic case



Ultra-relativistic case $\rightarrow a_0 \gg 1$

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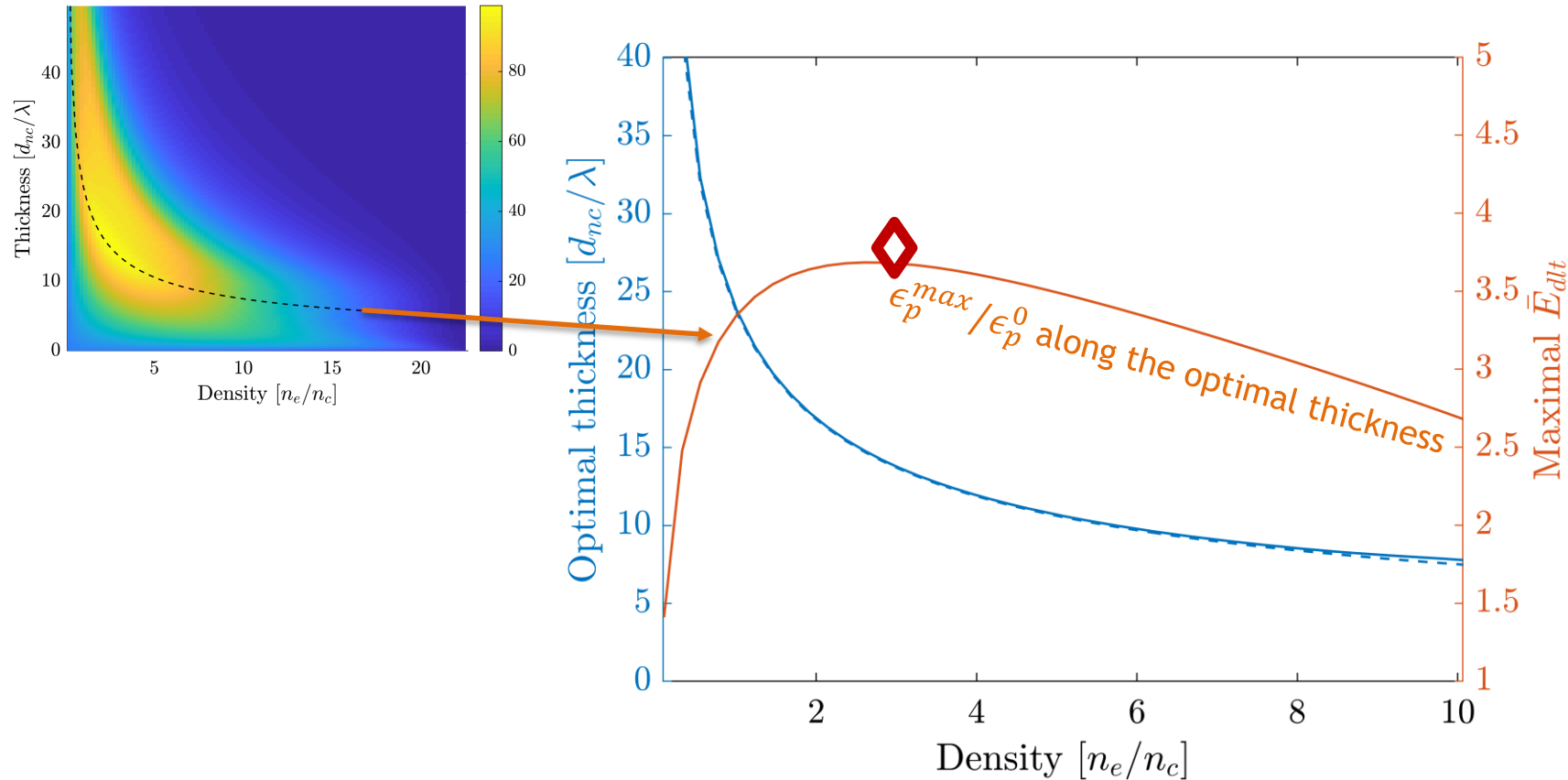
optimal density:

$$\bar{n}^{opt} = \frac{\lambda^2}{\pi w_0^2} \left(\frac{12\sqrt{2}\tau c/\lambda}{r_c^2 C_{nc}} \right)^{2/3}$$

Enhancement factor:

$$\bar{E}_{DLT}^{opt} = \frac{3C_{nc}}{4C_s} \left[1 - \frac{1}{\pi} \left(\frac{\sqrt{3}C_{nc}r_c^2}{2\tau c/\lambda} \right)^{2/3} \right]$$

Analytical solution in the ultra-relativistic case



Ultra-relativistic case $\rightarrow a_0 \gg 1$

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enhancement factors higher than ones reported in the literature

Conclusions

- Advanced TNSA via near-critical double-layer target

Acceleration with the Double-Layer Target (DLT)

The diagram compares two target configurations for TNSA acceleration:

- Flat solid foil:** Shows a laser pulse (red and blue) hitting a flat foil, resulting in a "hot electron cloud" (red and blue) that accelerates ions (blue).
- Near critical density layer:** Shows a laser pulse hitting a target with a near-critical density layer. The resulting electron cloud is "hotter and bigger" ($T_{e,DLT} > T_e$ and $n_{h,DLT} > n_h$). This leads to "Higher ions energy & number" and a higher maximum proton energy, $\epsilon_p^{max} \sim T_e \left[\log \left(\frac{n_h}{\bar{n}} \right) - 1 \right]$.

Advanced acceleration via near-critical DLT

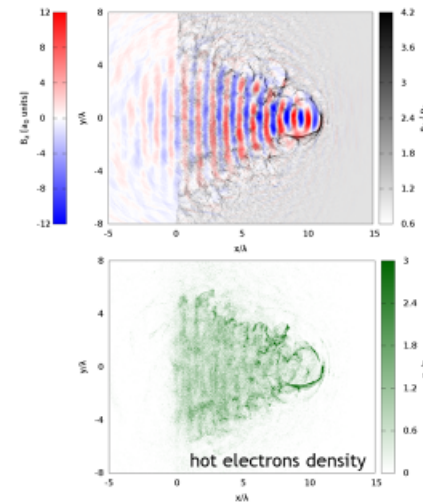
1 | Andrea Pazzaglia | CHILI

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- Modelization of proton acceleration with DLT

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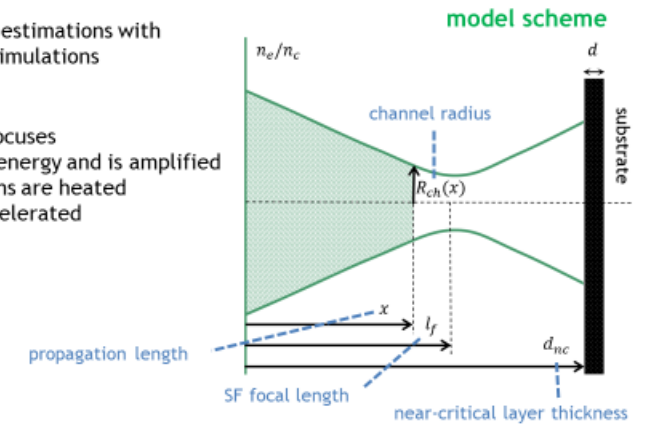


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MODEL STEPS:

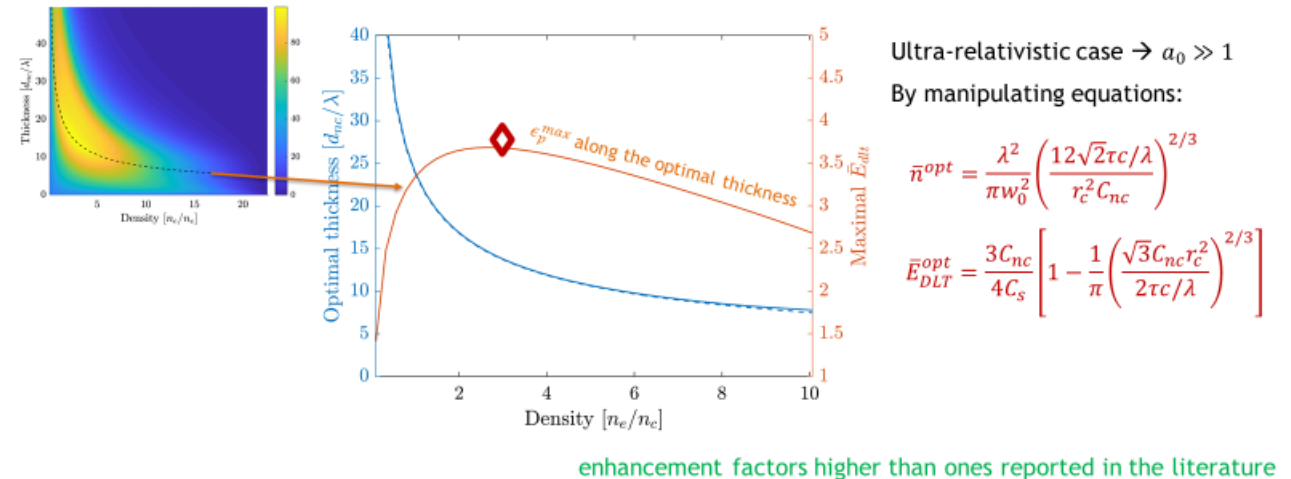
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2. Laser loses energy and is amplified
3. Hot electrons are heated
4. Ions are accelerated



Conclusions

- Advanced TNSA via near-critical double-layer target
- Modelization of proton acceleration with DLT
- Optimal DLT parameters

Analytical solution in the ultra-relativistic case



Future perspectives

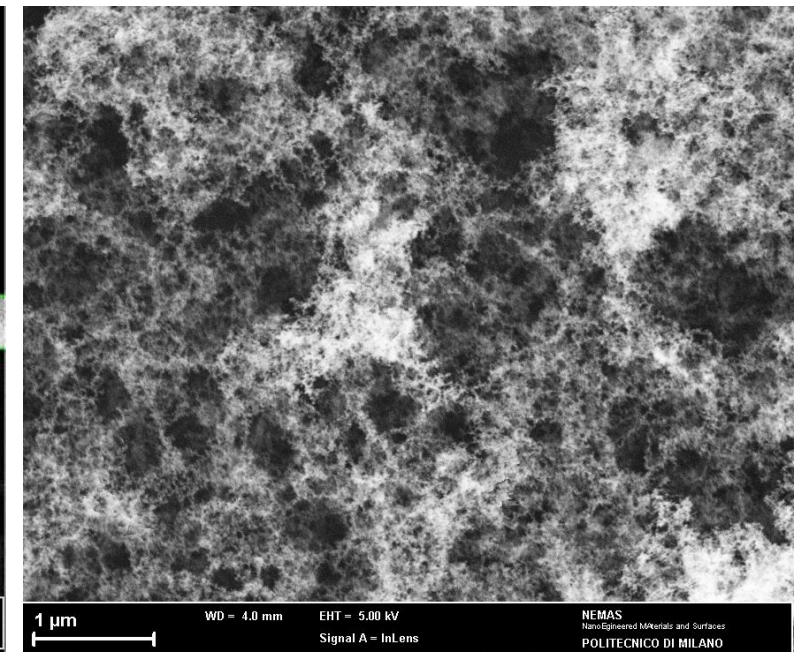
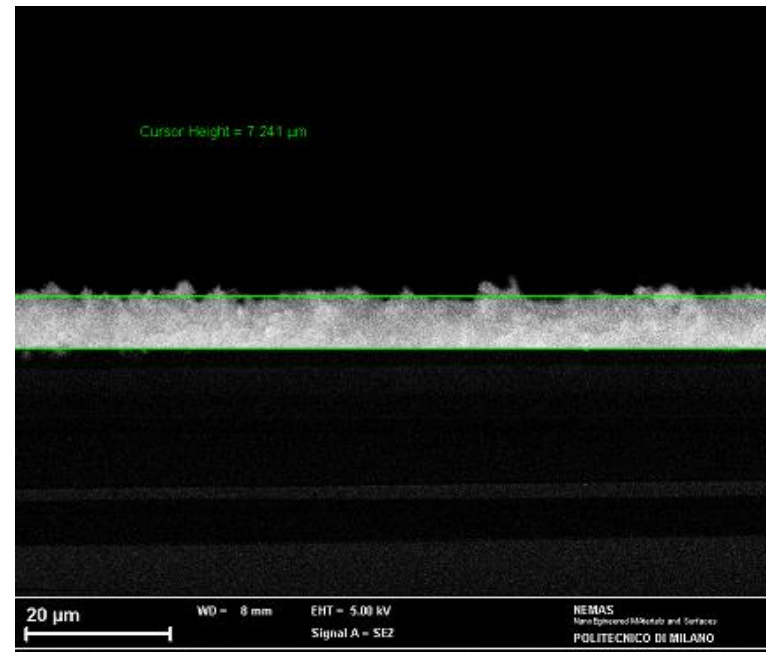
- Optimal DLT realization
 - ❖ Nanostructured near-critical layer

Zani, Alessandro, et al. *Carbon* 56 (2013): 358-365.

Maffini, A., et al. *Physical Review Materials* 3.8 (2019): 083404.

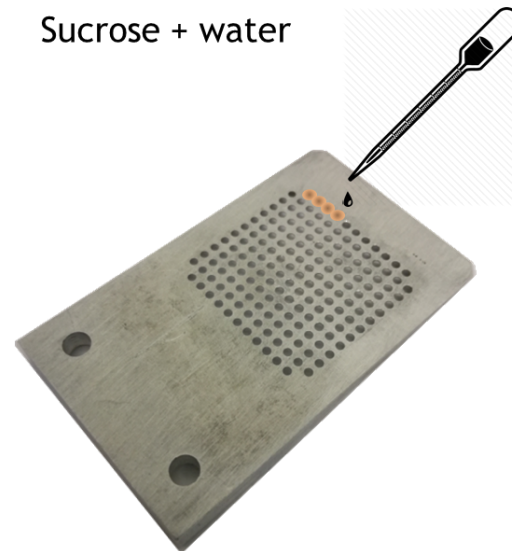
Passoni, Matteo, et al. *Plasma Physics and Controlled Fusion* (2019).

Produced by ns-PLD and fs-PLD

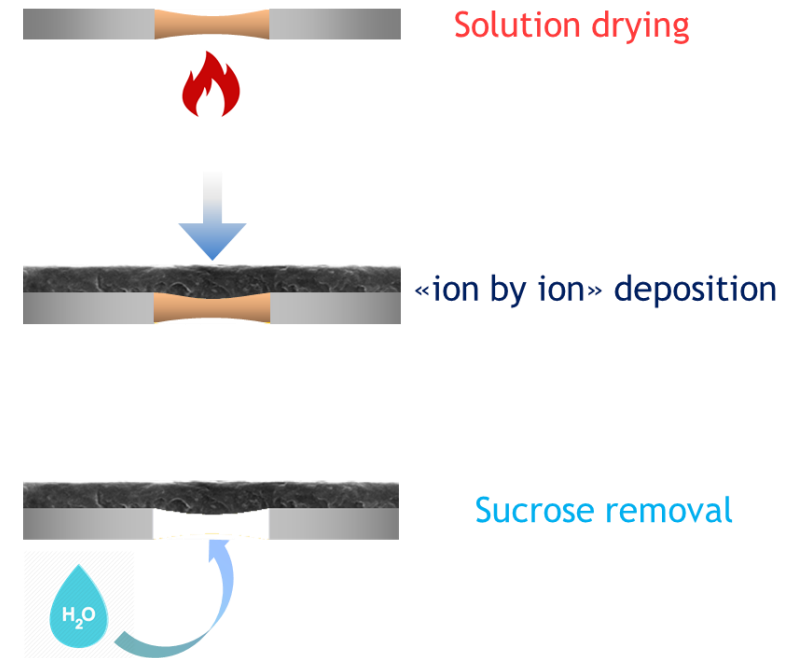
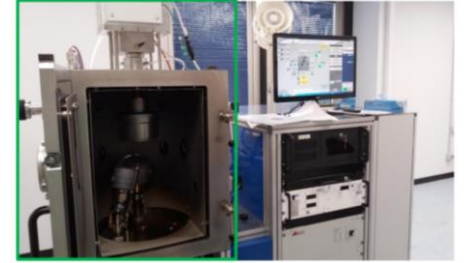


Future perspectives

- Optimal DLT realization
 - ❖ Nanostructured near-critical layer
 - ❖ Substrate production



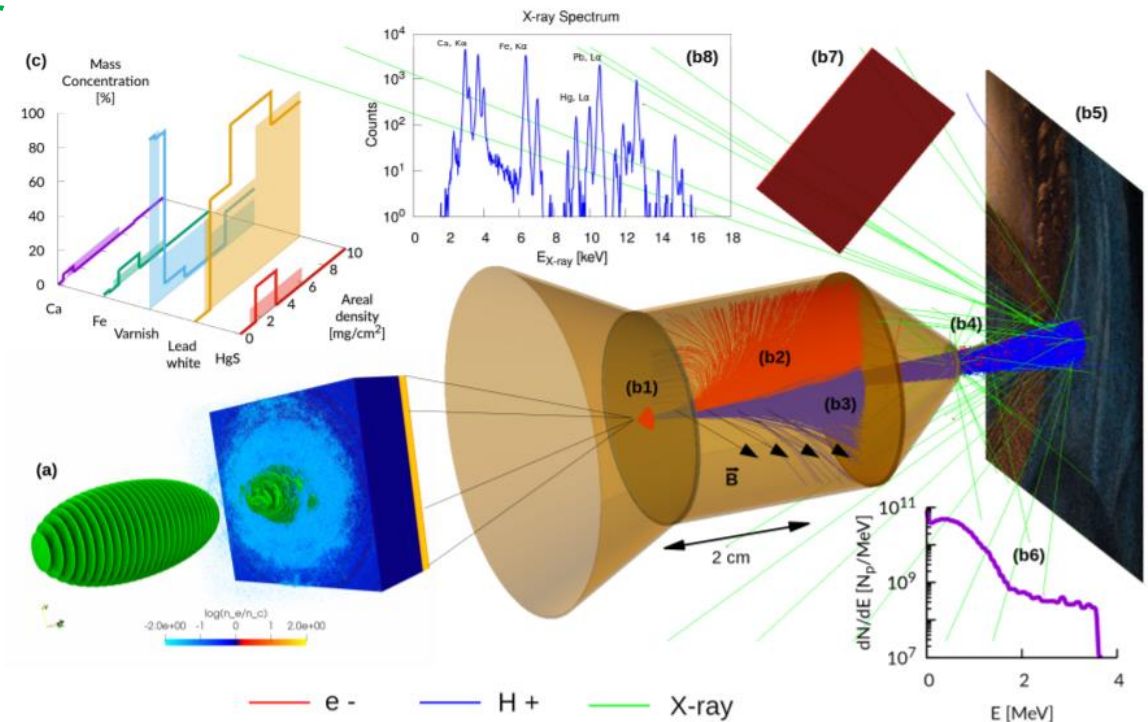
High Power Impulse
Magnetron Sputtering
(HiPIMS)



Future perspectives

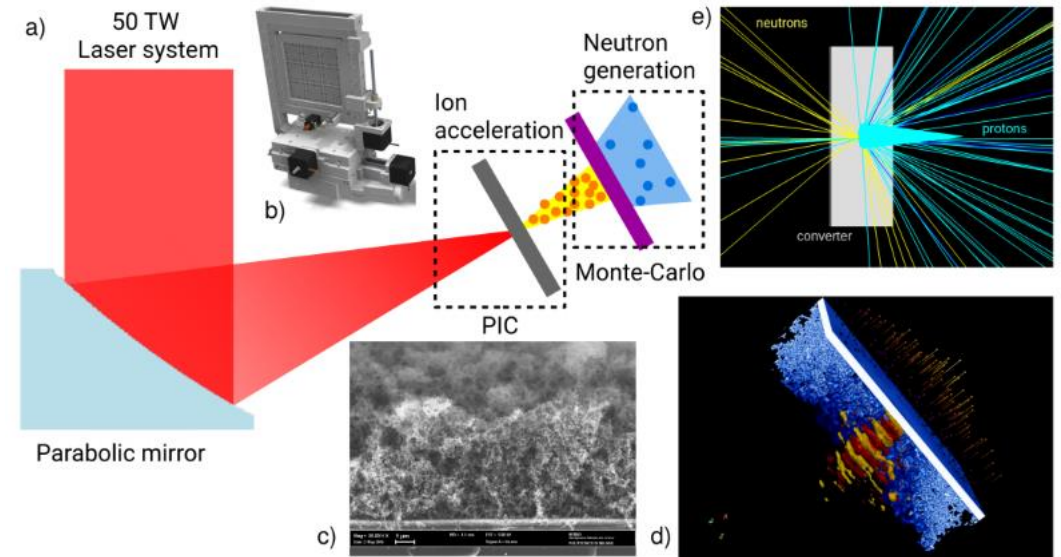
- Optimal DLT realization
 - ❖ Nanostructured near-critical layer
 - ❖ Substrate production
- Applications
 - ❖ Ion Beam Analysis

Passoni M., Fedeli L and Mirani F. Superintense Laser-driven Ion Beam Analysis (2019). *Scientific Reports*



Future perspectives

- Optimal DLT realization
 - ❖ Nanostructured near-critical layer
 - ❖ Substrate production
- Applications
 - ❖ Ion Beam Analysis
 - ❖ Neutron & Radioisotopes production



Fedeli L. et al. *New Journal of Physics* Under review

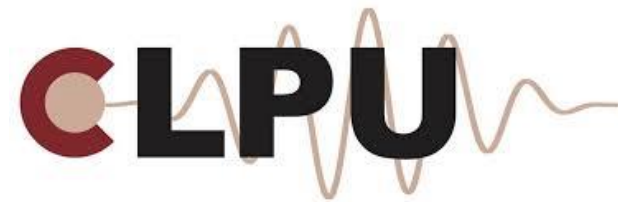
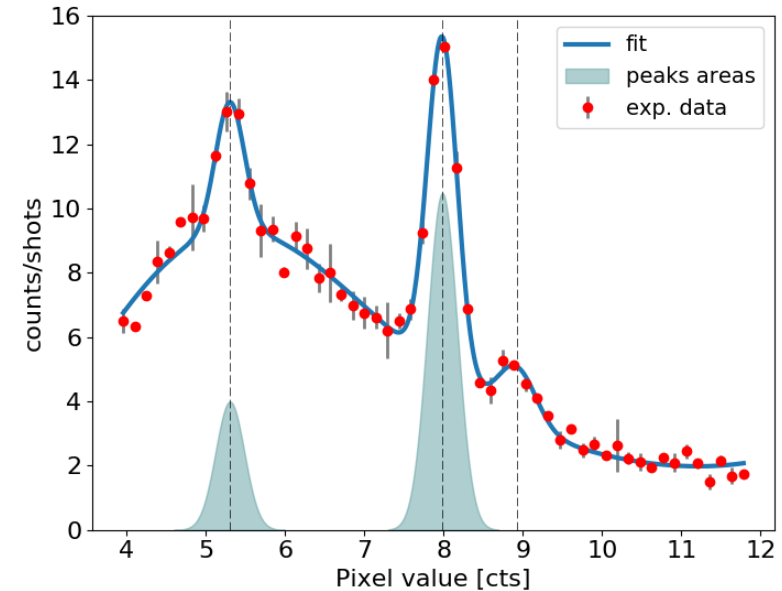
A. Tentori Master's thesis, Politecnico di Milano, Italy (2018)

F. Arioli Master's thesis, Politecnico di Milano, Italy (2019)

A. Giovannelli Master's thesis, Politecnico di Milano, Italy (2019)

Future perspectives

- Optimal DLT realization
 - ❖ Nanostructured near-critical layer
 - ❖ Substrate production
- Applications
 - ❖ Ion Beam Analysis
 - ❖ Neutron & Radioisotopes production
 - ❖ Dedicated experiments



Acknowledgments



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MILANO 1863



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ENSURE



M. Passoni



V. Russo



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D. Dellasega



A. Maffini



L. Fedeli



A. Pola



A. Formenti

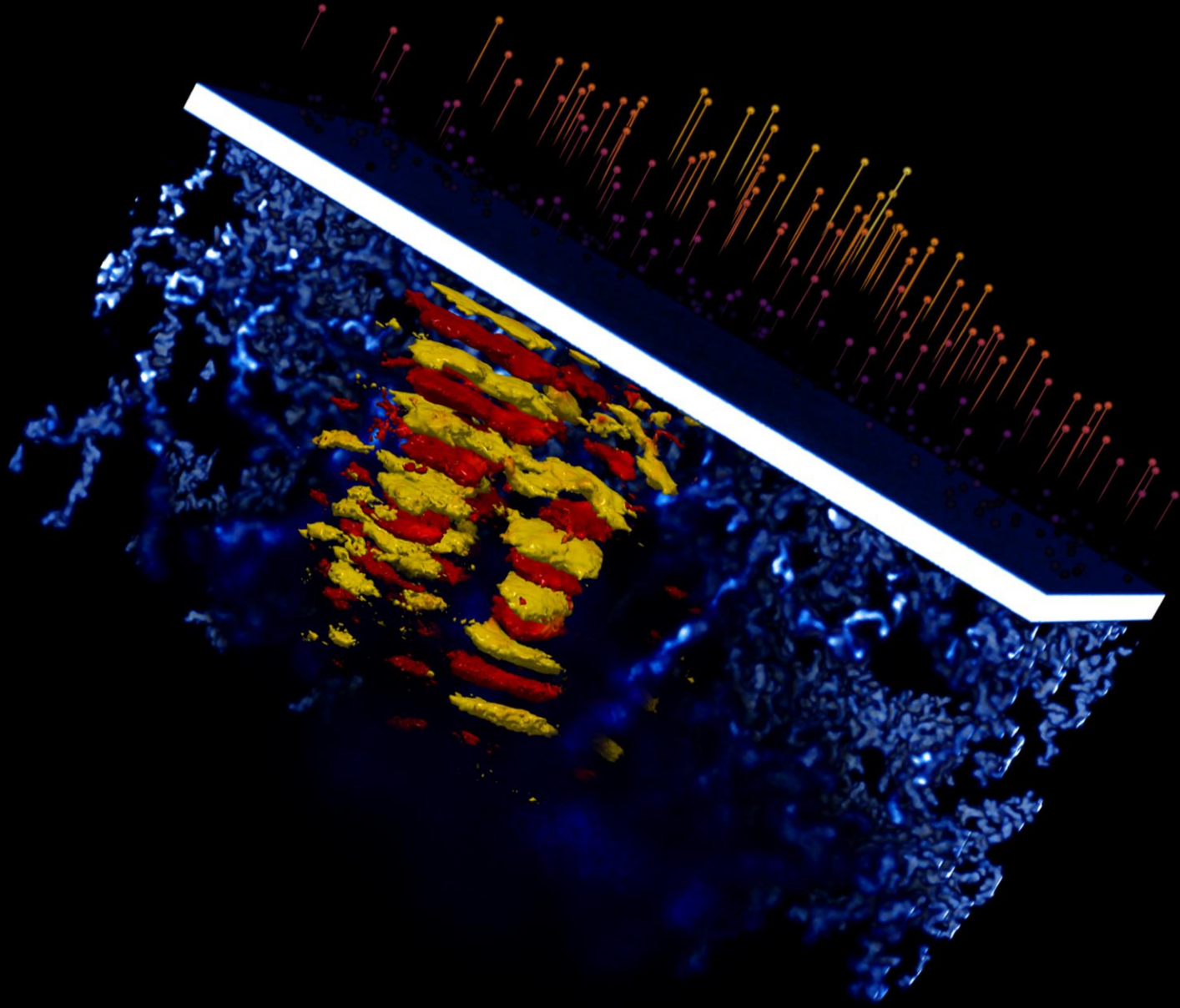


A. Pazzaglia



F. Mirani

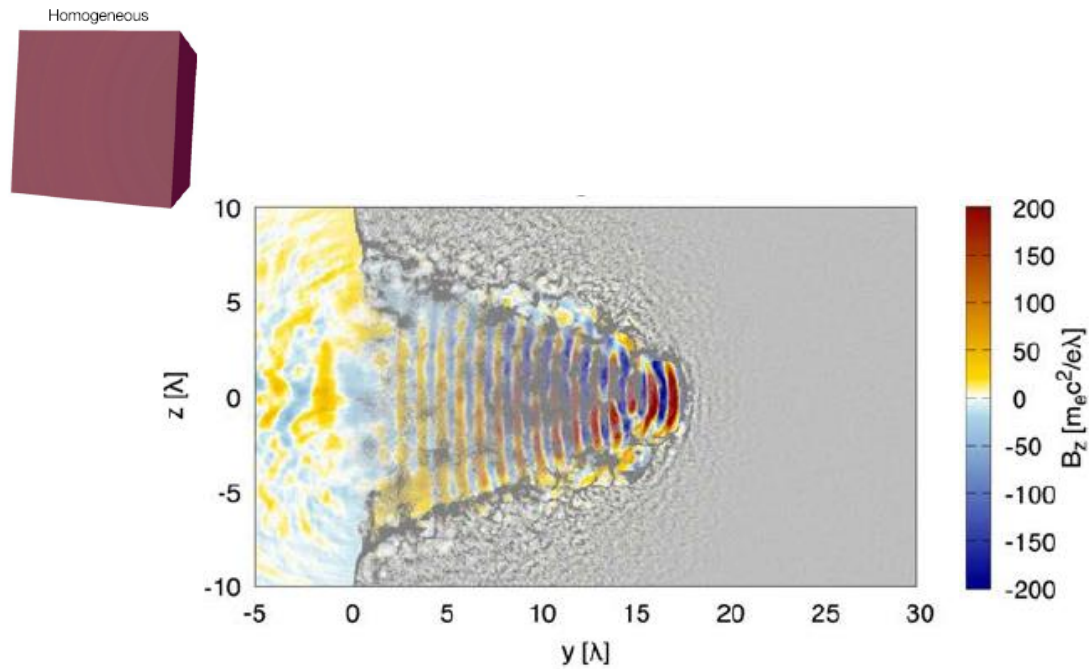




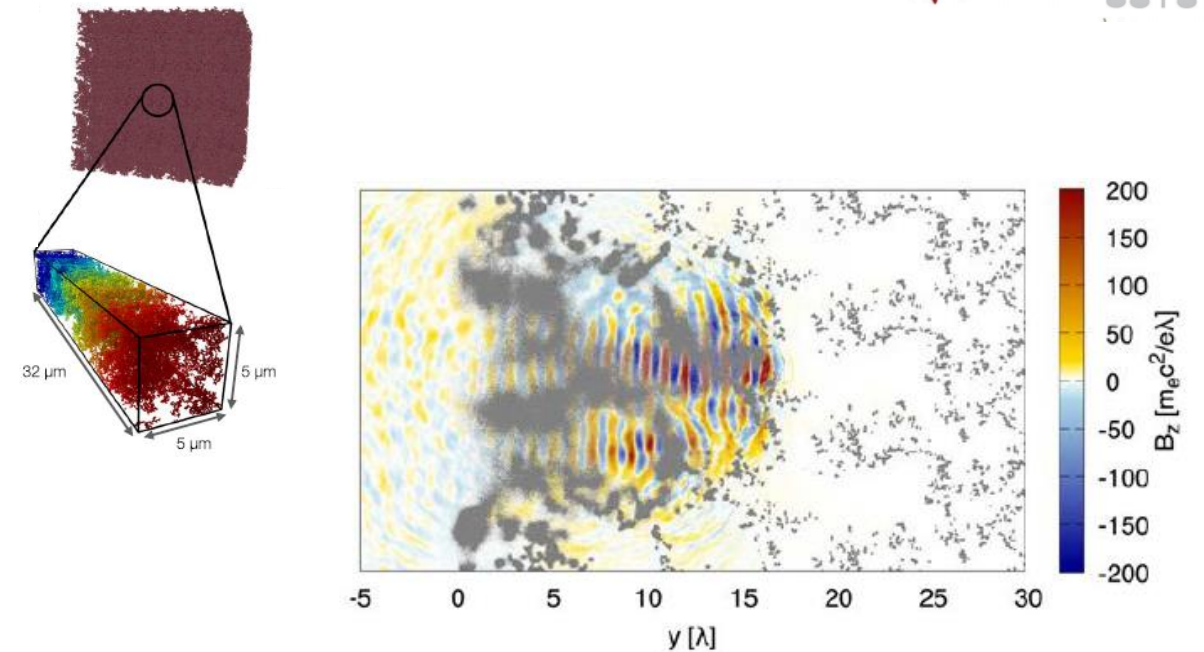
...and thank you for your attention!!

Realistic near-critical layer effects

Homogeneous near-critical plasma



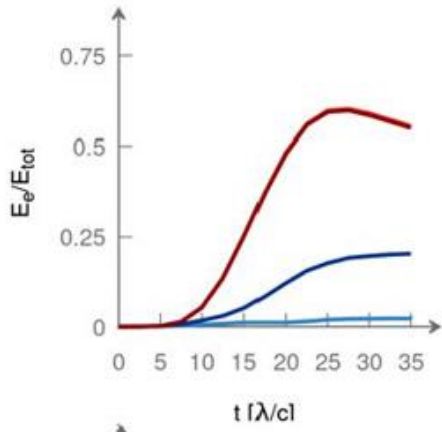
Nanostructured near-critical plasma



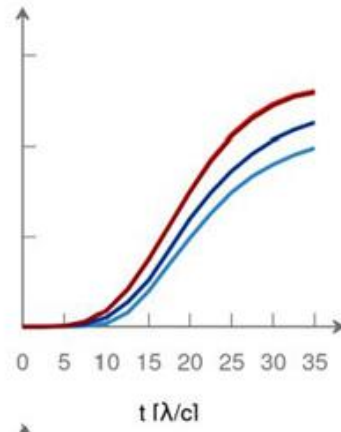
3D PIC simulations show evident differences

Differences in pulse energy loss and electron heating

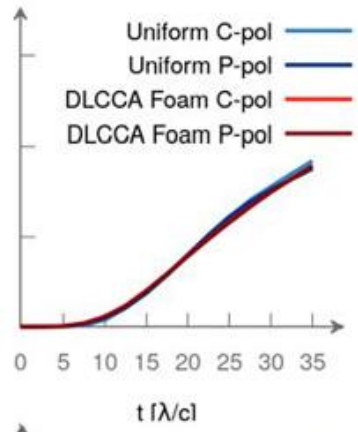
$a_0 = 5$
 $\bar{n} = 0.8$



$a_0 = 15$
 $\bar{n} = 0.3$



$a_0 = 45$
 $\bar{n} = 0.1$



If $\bar{n} \ll 1$ the energy loss is similar

the model previsions are valid in this regime!

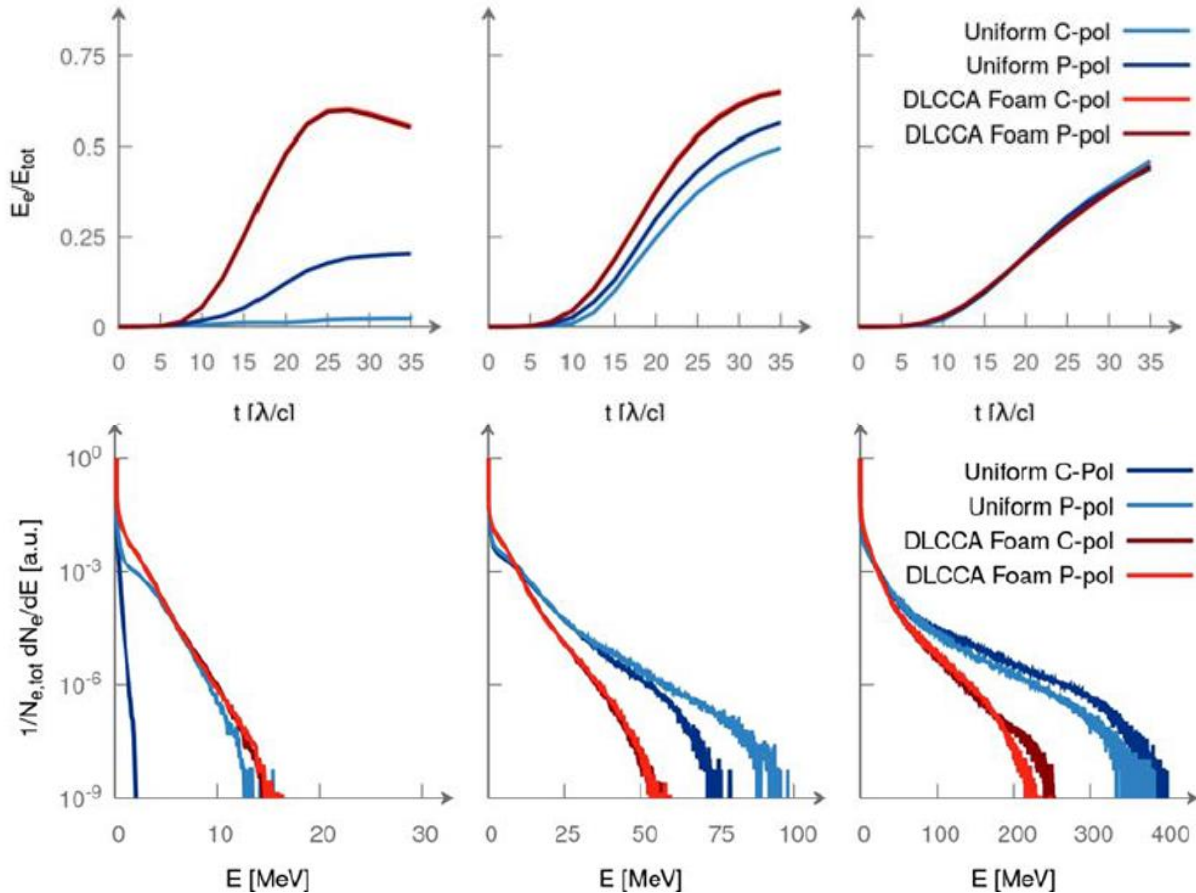
Differences in pulse energy loss and electron heating

$a_0 = 5$
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$a_0 = 15$
 $\bar{n} = 0.3$

$a_0 = 45$
 $\bar{n} = 0.1$

If $\bar{n} \ll 1$ the energy loss is similar



the model previsions are valid in this regime!

Nanostructured plasma:

- Electrons temperature ↓
- Electrons number ↑

$$\epsilon_p^{max} = T_{DLT} \left[\log \left(\frac{n_{h,DLT}}{\tilde{n}} \right) - 1 \right] \rightarrow \text{lower energy expected}$$

uniform near-critical plasma are optimal